



Structural defects in SiO₂/SiC interface probed by a slow positron beam

M. Maekawa^{a,*}, A. Kawasuso^a, Z.Q. Chen^a, M. Yoshikawa^b,
R. Suzuki^c, T. Ohdaira^c

^aAdvanced Science Research Center, Japan Atomic Research Energy Institute, 1233 Watanuki, Takasaki, Gunma 370-1292, Japan

^bTakasaki Establishment, Japan Atomic Research Energy Institute, Takasaki, Japan

^cPhotonics Research Institute, National Institute of Advanced Industrial Science and Technology, Takasaki, Japan

Received 31 May 2004; accepted 1 October 2004

Abstract

We studied structure of SiO₂/SiC interfaces grown by thermal oxidization using low energy positron beams. From the positron lifetime measurements, typical two lifetime components related to amorphous network and large open spaces are observed in SiO₂, while only bulk lifetime in SiC. The lifetime of positrons near the interface is similar to that of SiO₂ amorphous network and no lifetime components related to large open spaces are detected. From the Doppler broadening measurements, the momentum distribution, probably related to oxygen valence electrons, is enhanced near the interface. The above results imply that the interface region has less open spaces than SiO₂ involving many oxygen dangling bonds.

© 2005 Elsevier B.V. All rights reserved.

PACS: 68.35.C; 78.70.B

Keywords: SiO₂/SiC interface; Interface defects; Positron annihilation spectroscopy; Dry oxidation; Pyrogenic oxidation

1. Introduction

Silicon carbide (SiC) is a wideband gap semiconductor which is expected as high power device applications. Above all, 4H-SiC has an advantage to high power devices due to its wide bandgap (3.2 eV) [1]. One of the most attractive features of SiC is that

the oxide layers needed for metal-oxide-semiconductor (MOS) devices can be grown by conventional thermal oxidation. However, it is reported that many interfacial defects remain near the SiO₂/SiC interface and these degrade performance of the MOS-based field effect transistor (MOSFET). Ordinarily produced SiO₂/SiC interface still remains in the range of 10¹¹ to 10¹² cm⁻², which is about two orders of magnitude higher than those for SiO₂/Si interfaces [2–4]. The reduction of these centers is an important issue to realize of the high performance MOS device. In recent

* Corresponding author. Tel.: +81 27 346 9330;
fax: +81 27 346 9324.

E-mail address: maekawa@taka.jaeri.go.jp (M. Maekawa).

years, although the pyrogenic oxidation improve the electrical properties of interface as compared to dry oxidation [5], the detailed mechanism is not clarified. The structure of SiO₂/SiC interfaces should be revealed.

Positron annihilation spectroscopy (PAS) is a powerful tool to detect open volume defects in solids [6,7]. Using slow positron beam, the depth distribution of defects near subsurface region of several microns can be obtained. The PAS technique can be also used to characterize defects near SiO₂/SiC interfaces. Indeed, previous positron annihilation studies suggest the existence of open volume type defects at SiO₂/Si interfaces and their passivation with some specific treatments [8–10]. In this study, we investigate the characteristics of the SiO₂/SiC interface in SiC–MOS grown pyrogenic and dry oxidation methods using PAS.

2. Experimental

The samples used in this study (10 mm × 10 mm) were cut from CVD-grown n-type 4H–SiC epilayers on 8°-off oriented 4H–SiC(0 0 0 1) substrates purchased from CREE Research Inc. These were subjected to the pyrogenic oxidation at 1100 °C for 2 h or to the dry oxidation at 1200 °C for 3 h. The oxide thickness was approximately 40 nm. Slow positrons generated by a ²²Na source with the energies of 0.2–20 keV are injected into specimens. Implanted positrons annihilate with electrons to emit two gamma quanta (511 keV) per one positron–electron pair. The energy spectrum is Doppler broadened depending on electron momenta. The Doppler spectrum was measured by a high energy-resolution Ge detector. *S* and *W* parameters defined as the normalized peak and tail intensity of Doppler spectrum is determined [8]. In this study, the peak region for *S* parameter is $p = (0–2) \times 10^{-3}m_{0c}$ and the tail region for *W* parameter is $p = (10–15) \times 10^{-3}m_{0c}$. To obtain the chemical surroundings of positron annihilation sites, the Doppler spectrum with enhanced signal-to-noise ratio are measured using the coincidence Doppler broadening (CDB) technique [11]. In this study, p-type SiC was used as a reference because of negligible positron trapping centers inside it. Moreover, positron lifetime measurements were also carried out using the

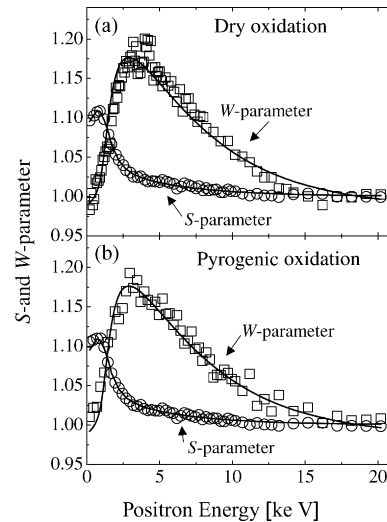


Fig. 1. *S* and *W* parameters as a function of incident positron energy for SiO₂/SiC structures fabricated by dry oxidation (a) and pyrogenic oxidation (b). Open circles denote *S* parameters and open rectangular are *W* parameters. These parameters are normalized to the bulk values. Lines are fitting result obtained from VEPFIT code.

AIST pulsed positron beam. After subtraction of the background components, the lifetime spectra were analyzed using the PATFIT program.

3. Results and discussion

Fig. 1 shows *S* and *W* parameters as a function of incident positron energy obtained after dry and pyrogenic oxidation process. In order that *S* and *W* parameter is influenced by the annihilation ratio with valence and core electrons respectively, change of these parameters is usually contrary to each other. *S* and *W* parameters remain constant at $E > 15$ keV. This reflects the annihilation of positrons in the SiC epitaxial layers. At $E < 1$ keV, most positrons stop in the SiO₂ layers. The enhanced *S* parameter is caused by the small energy distribution of positronium formed in the SiO₂ amorphous network. The small *W* parameter in this region is also reflecting positronium annihilation. A little reduction of *S* parameter at $E = 0.2$ keV indicates the positron annihilation at the surface. The energy range corresponding to the SiO₂/SiC interface is $E = 3–6$ keV. Here, no significant change of *S* parameter is found. On the contrary, *W* parameter dramatically

changes in the same energy range. This is a remarkable phenomenon. Under an assumption that positrons annihilate in the ideal interface of SiC and SiO₂, S parameter is expressed with the linear combination of S parameters of each layer by giving weight with an annihilation ratio x , like $S = xS_{\text{SiO}_2} + (1-x)S_{\text{SiC}}$. However, in this region, change of S and W parameter cannot be expressed with the same change of x because W parameter changes with an extremum although S parameter changes in monotone. This suggests that the existence of a transition region that is not a simple mixture of SiO₂ and SiC but another chemical state with different electron momentum distribution is detected. For the above two types of SiO₂/SiC samples, both S - E and W - E plots are nearly the same.

To determine the electron momentum distribution more in detail, CDB measurements were performed. Ratio curves to the SiC bulk are shown in Fig. 2. The intensity for SiO₂ region is enhanced at $p < 3 \times 10^{-3} m_0c$. This corresponds to the high S parameter in Fig. 1 because of the positronium formation. A large bump also appears at $p = (10\text{--}20) \times 10^{-3} m_0c$. In the case of SiO₂/SiC transition region, the intensity at $p < 3 \times 10^{-3} m_0c$ is rather reduced, while that at $p = 10\text{--}20 \times 10^{-3} m_0c$ is fairly pronounced. This enhancement is not simply explained as a relative increase due to the balance for the decrease at $p < 3 \times 10^{-3} m_0c$. Indeed, Fig. 1 shows no significant decrease in S parameter where W parameter increases ($E \sim 3$ keV). The bump at $p = (10\text{--}20) \times 10^{-3} m_0c$ may be explained as (i) the missing of the neck for the second Brillouin zone which clearly appears in the case of crystalline structure, i.e., in SiC region, or (ii)

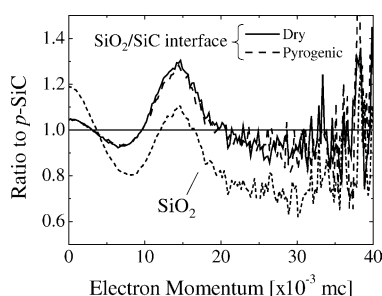


Fig. 2. Ratio curves of electron momentum distribution for the SiO₂ and interface to the SiC bulk observed by coincidence Doppler broadening measurement.

Table 1
Parameters obtained by fitting the positron lifetime spectrum of SiO₂ and interface layer

	SiO ₂ ($E = 0.8$ keV)		Interface ($E = 3.7$ keV)
	τ_1 (ns)	τ_2 (ns)	τ_1 (ns)
Dry	0.408 (39.3%)	1.63 (60.7%)	0.451 (100%)
Pyrogenic	0.419 (41.0%)	1.58 (59%)	0.470 (100%)

the chemical environment of positron annihilation sites is different from that in SiC matrix. However, since the neck of the second Brillouin zone of SiC is located at $p = (8\text{--}15) \times 10^{-3}$ [12], the former possibility is not favored. It is known that positrons are attracted to oxygen atoms rather strongly [13–16]. Thus, as for the second possibility, we may consider that positrons preferentially annihilate with valence electrons of oxygen atoms.

Table 1 shows the results of positron lifetime measurements for SiO₂ layer and SiO₂/SiC interface. In the SiO₂ layer, two lifetime components were determined; one is about 400 ps attributed to free annihilation in SiO₂ amorphous network and the other 1.6 ns to pick-off annihilation of ortho-positronium. This means that in the SiO₂ layer there are large voids to allow the positronium formation. In the interface region, only a single positron lifetime component of about 450 ps is resolved. This result shows that the structure of interface layer is rather porous but the positronium formation is forbidden. Probably, this is explained as the reduced concentration of large voids at the interface. Alternatively, though the interface resembles with SiO₂, a large amount of defects exists to reduce the positronium formation as competitive trapping centers.

From the above results, the structure of the interface transition region is not simple admixture between SiO₂ and SiC. By the way, there are almost no differences of PAS result between specimens obtained by dry and pyrogenic oxidation. However, from the C - V measurements the net number of interface traps per unit area (N_{it}) obtained for SiO₂/SiC structure by the dry oxidation is one order of magnitude higher than that made by the pyrogenic oxidation [17]. This is considered that electrically active centers may be terminated by steam (moisture existing oxidation ambient). On the contrary, the oxygen related defects

are observed in PAS measurements. Those may be attributed to the imperfect oxides generated during oxidation process because interface structure is not simple transition layer from SiO₂ to SiC. This shows that the moisture existing oxidation ambient works the reduction of the interface state density, but have little effect on the formation of the imperfect oxides. Such defects are formed due to the lack of oxygen in oxidization. It is suggested that a reoxidation process is valid for the improvement of the interface properties. For the clarification of interface properties more in detail, it is necessary to measure SiO₂/SiC interfaces fabricated with the technique of reducing interface state density, such as post-oxidation annealing.

4. Summary

We studied structural defects at SiO₂/SiC interface grown by dry and pyrogenic oxidization using slow positron beams. From the Doppler broadening measurement, *W* parameter dramatically increases in the energy range corresponding to the SiO₂/SiC interface. The CDB measurements show that the chemical environment of positron annihilation sites in transition layer is different from that in SiC matrix, and show that positrons annihilate with valence electrons of oxygen atoms due to its strong affinity for positrons. From the positron lifetime measurements, the structure of interface layer is porous without large voids to allow the positronium formation. As a result, oxygen related defects exist in the interface transition layer. Those defects are probably concerned with the imperfect oxides because of its SiO₂-like porous structure. There are almost no

differences of PAS result between specimens obtained by dry and pyrogenic oxidation. This shows that no significant contributions of the moisture in oxidization ambient to the formation of such defects.

References

- [1] G. Harris, Properties of Silicon Carbide, The Institute of Electrical Engineers, London, 1995.
- [2] V.V. Afanas'ev, M. Bassler, G. Pensl, A. Stesmans, Mater. Sci. Forum 389–393 (2002) 961.
- [3] K. Fukuda, J. Senzaki, M. Kushibe, K. Kojima, R. Kosugi, S. Suzuki, S. Harada, T. Suzuki, T. Tanaka, K. Arai, Mater. Sci. Forum 389–393 (2002) 1057.
- [4] T. Iida, Y. Tomioka, H. Yaguchi, M. Yoshikawa, Y. Ishida, H. Okumura, S. Yoshida, Jpn. J. Appl. Phys. 39 (2000) L1054.
- [5] M. Yoshikawa, M. Satoh, T. Oshima, H. Itoh, Mater. Sci. Forum 389–393 (2002) 1009.
- [6] P.J. Schultz, K.G. Lynn, Rev. Mod. Phys. 60 (1988) 701.
- [7] R. Krause-Rehberg, H.S. Leipner, Positron annihilation in semiconductors, Springer Series in Solid-State Sciences, vol. 27, Springer, Berlin, 1998.
- [8] A. Uedono, S. Tanigawa, Y. Ohji, Phys. Lett. A133 (1988) 82.
- [9] P. Asoka-Kumar, K.G. Lynn, D.O. Weich, J. Appl. Phys. 76 (1994) 4935.
- [10] T.C. Leung, P. Asoka-Kumar, B. Nielsen, K.G. Lynn, J. Appl. Phys. 73 (1993) 168.
- [11] K.G. Lynn, J.R. MacDonald, R.A. Boie, L.C. Feldman, J.D. Gabbe, M.F. Fobbins, E. Bonderup, J. Golovchenko, Phys. Rev. Lett. 38 (1977) 241.
- [12] Z. Tang, T. Nonaka, Y. Nagai, M. Hasegawa, Mater. Sci. Forum 363–365 (2001) 67.
- [13] U. Myler, Appl. Phys. Lett. 68 (1996) 3333.
- [14] N.Yu. Arutyunov, R. Krause-Rehberg, Solid State Phenomena 95–96 (2003) 507.
- [15] K. Ito, Y. Kobayashi, Appl. Phys. Lett. 82 (2003) 654.
- [16] Y. Nagai, T. Nonaka, M. Hasegawa, Y. Kobayashi, C.L. Wang, W. Zheng, C. Zhang, Phys. Rev. B 60 (1999) 11863.
- [17] M. Maekawa, A. Kawasuso, M. Yoshikawa, H. Itoh, Appl. Surf. Sci. 216 (2003) 365.