



Development of Monte-Carlo Scheme for Runaway Electron Generation and Confinement Code

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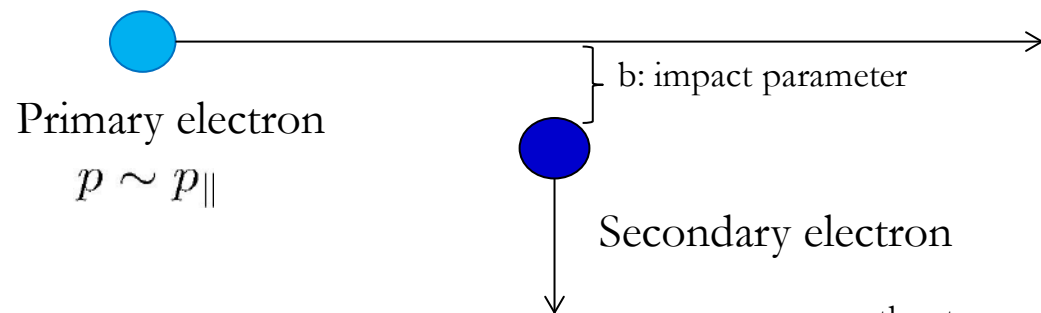
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Runaway Electron is a big issue in ITER ...

- Runaway Electrons generated during disruption may cause substantial damage to plasma facing components. **In ITER, significant runaway generation is predicted theoretically,** for which avalanche amplification plays an essential role.

Avalanche amplification due to close collisions

(Rosenbluth & Putvinski, 1997)



Current Amplification Rate

$$\gamma \tau_{CQ} \sim \frac{I_p}{I_A \ln \Lambda}$$

γ : growth rate

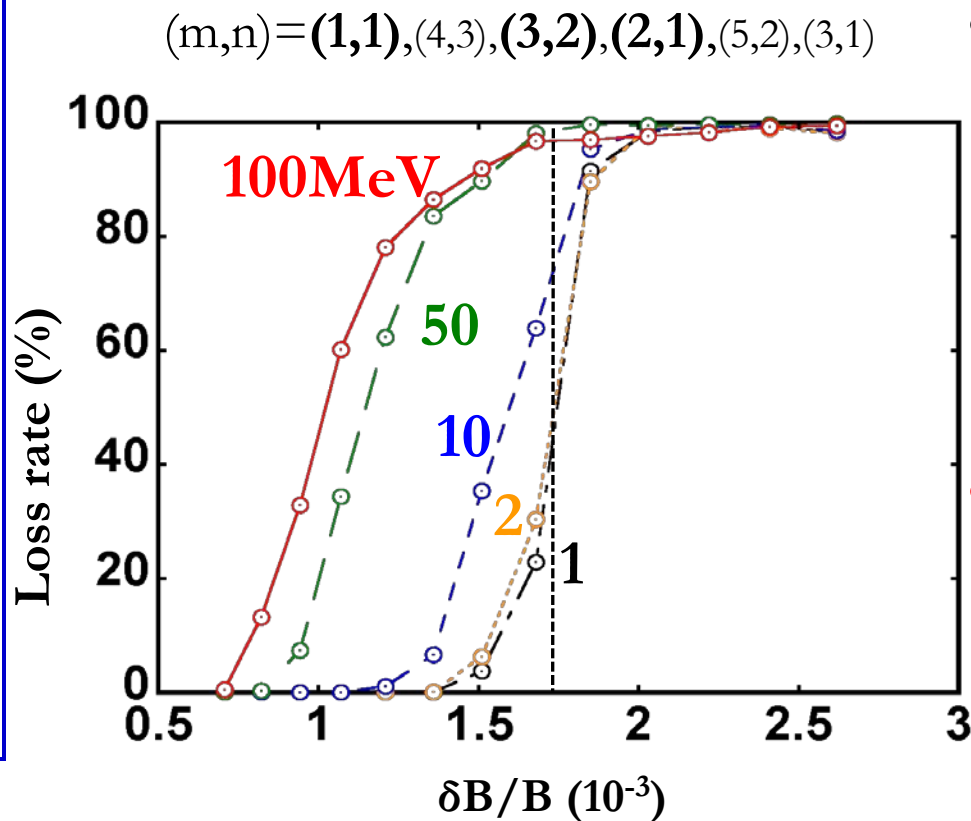
τ_{CQ} : current quench time

I_p : plasma current

I_A : Alfvén current

ETC-Rel code

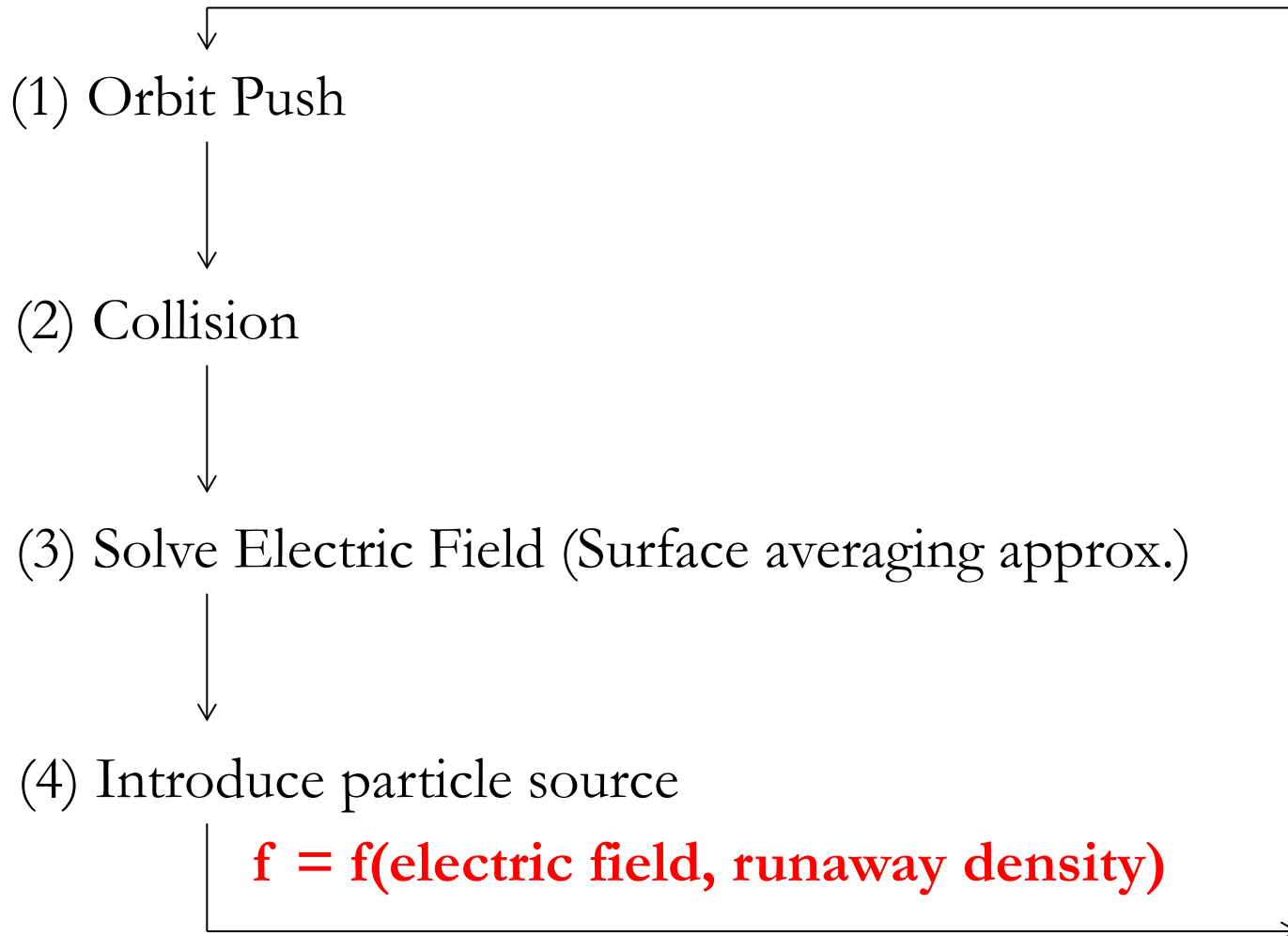
- Collisionless orbit simulation ETC-Rel has been developed in JAEA (Tokuda 1999-; Matsuyama, et al., 2012-)



- Orbit loss rate in presence of low-order magnetic perturbation has been studied previously (Matsuyama, et al., submitted to JPSJ suppl.)
- This work extends ETC-Rel to include runaway generation mechanisms with self-consistent electric field model.

New version of ETC-Rel

- Self-consistent runaway source and electric field model is newly implemented in ETC-Rel.



Orbit push

- Original version of ETC-Rel code solves collisionless orbit of runaway electrons for given energy and pitch angle with the relativistic guiding-center equations.

- Guiding-center equations

(Cary & Brizard, 2000)

$$\frac{d\mathbf{X}}{dt} = \frac{p_{\parallel}}{m\gamma} \frac{\mathbf{B}^{**}}{B_{\parallel}^{**}} + \mathbf{E}^* \times \frac{\mathbf{b}}{B_{\parallel}^{**}}$$

$$\frac{dp_{\parallel}}{dt} = e \frac{\mathbf{B}^{**}}{B_{\parallel}^{**}} \cdot \mathbf{E}^*$$

$$\mathbf{E}^* = \mathbf{E} - \frac{\mu}{e} \nabla B - \frac{p_{\parallel}}{e} \frac{\partial \mathbf{b}}{\partial t}$$

$$\mathbf{B}^* = \mathbf{B} + \frac{p_{\parallel}}{e} \nabla \times \mathbf{b}$$

$$\mathbf{B}^{**} = \mathbf{B}^* + \delta \mathbf{B} \quad \delta \mathbf{B} = \nabla \times (\alpha \mathbf{B})$$

- Interface to MEUDAS equilibrium code for realistic tokamak geometry
- MPI-OpenMP hybrid code (tested 32-128 node in Helios)
- Solving GC eqs. in Cylindrical coordinates including the vacuum region.
- Possible to take into account magnetic perturbations into the calculation.

Relativistic Collisions

- Monte-Carlo Collision Operator [G. Papp, et al., (2011)]

(1) Pitch angle scattering

$$\frac{d}{dt}\xi = -\mathcal{I}(q)\xi \quad \frac{d}{dt}\sigma_\xi^2 = \mathcal{I}(q)(1 - \xi^2)$$

(2) Slowing-down and energy scattering

$$\frac{d}{dt}q = \frac{1}{\tau q^2} \left\{ -\mathcal{J}(q)(1 + q^2) + \frac{\partial}{\partial q}[\mathcal{J}(q)\mathcal{P}(q)] \right\} \quad \frac{d}{dt}\sigma_q^2 = \frac{2}{\tau q^2} \mathcal{J}(q)\mathcal{P}(q)$$

$$\mathcal{J}(q) = \frac{q^2}{\epsilon(1 + q^2)} G\left(\frac{q}{\sqrt{2\epsilon(1 + q^2)}}\right) \quad \epsilon = v_t^2$$

$$\mathcal{P}(q) = \frac{\epsilon \sqrt{1 + q^2}^3}{q}$$

$$\mathcal{I}(q) = \frac{\sqrt{1 + q^2}}{\tau q^3} \left[Z_{\text{eff}} + \Phi\left(\frac{q}{\sqrt{2\epsilon(1 + q^2)}}\right) - G\left(\frac{q}{\sqrt{2\epsilon(1 + q^2)}}\right) + \epsilon \frac{q^2}{1 + q^2} \right]$$

Generation Source Model

- Source model

Primary electron generation (Dreicer acceleration, Connor & Hastie NF)

$$\dot{n}_r^I \simeq \frac{n_e}{\tau} \left(\frac{m_e c^2}{2T_e} \right)^{3/2} \left(\frac{E_D}{E_{\parallel}} \right)^{3(1+Z_{\text{eff}})/16} \exp \left(-\frac{E_D}{4E_{\parallel}} - \sqrt{\frac{(1+Z_{\text{eff}})E_D}{E_{\parallel}}} \right)$$

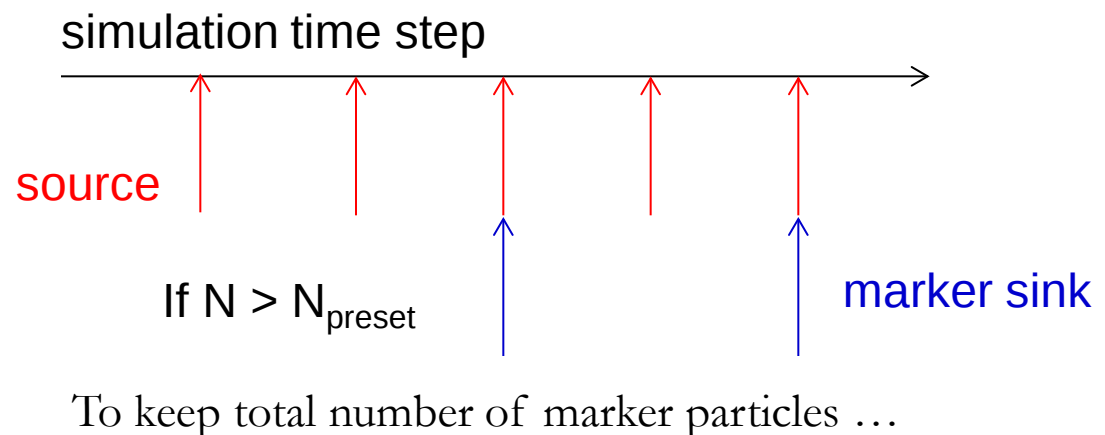
Secondary electron generation (Avalanche amplification, Rosenbluth & Putvinski NF)

$$\dot{n}_r^{II} \simeq n_r \frac{E_{\parallel}/E_c - 1}{\tau \ln L} \sqrt{\frac{\pi\varphi}{3(Z_{\text{eff}} + 5)}} \left(1 - \frac{E_c}{E_{\parallel}} + \frac{4\pi(Z_{\text{eff}} + 1)^2}{3\varphi(Z_{\text{eff}} + 5)(E_{\parallel}^2/E_c^2 + 4/\varphi^2 - 1)} \right)^{-1/2}$$

Dreicer field $E_D = \frac{m_e^2 c^3}{e\tau T_e}$

Critical field $E_c = \frac{m_e c}{e\tau}$

Collision times $\tau = \frac{4\pi\epsilon_0^2 m_e^2 c^3}{n_e e^4 \ln \Lambda}$



1D Electric Field Solver

- 1D model for solving parallel electric field

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial E_{\parallel}}{\partial r} \right) = \mu_0 \frac{\partial}{\partial t} (\sigma_{\parallel} E_{\parallel} + j_r)$$

- Space: Finite Difference (Staggered mesh)
- Time: Fully Implicit/Crank-Nicholson Scheme

- **Consider response of the runaway current to temperature drop**

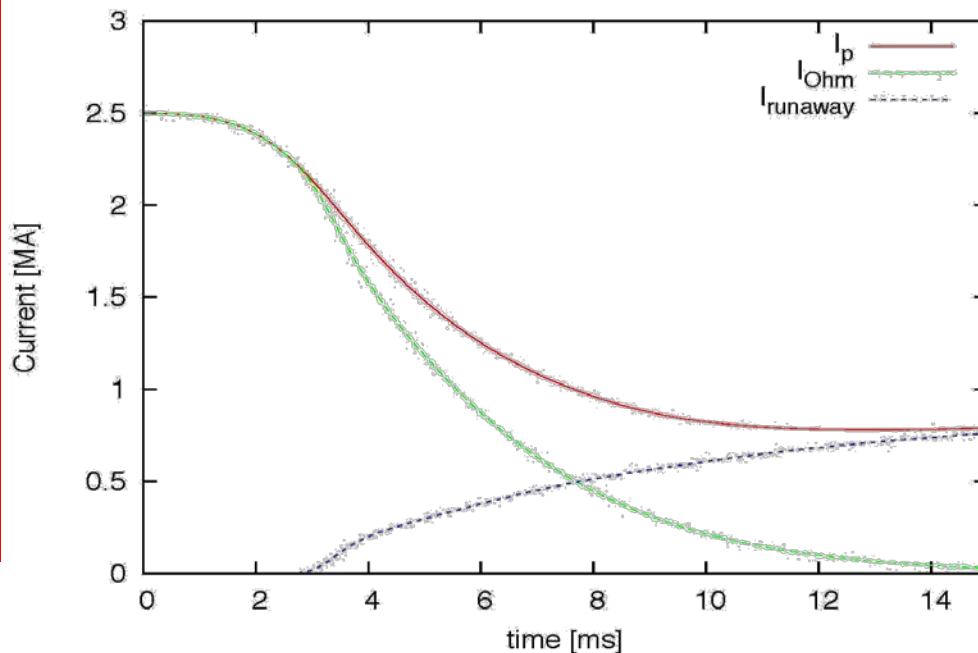
$$T_e(t, r) = T_{\text{final}}(r) + [T_0(r) - T_{\text{final}}(r)] e^{-t/t_0}$$

- Runaway density and current are evaluated from ETC-Rel code with surface averaging (**full-f approach**).

$$n_r(\rho_i) = \frac{1}{\Delta V} \sum_i w_s^i, \quad w_s = N_p / N_s, \quad N_p \equiv \int d\rho n_r(\rho) \frac{dV}{d\rho}$$

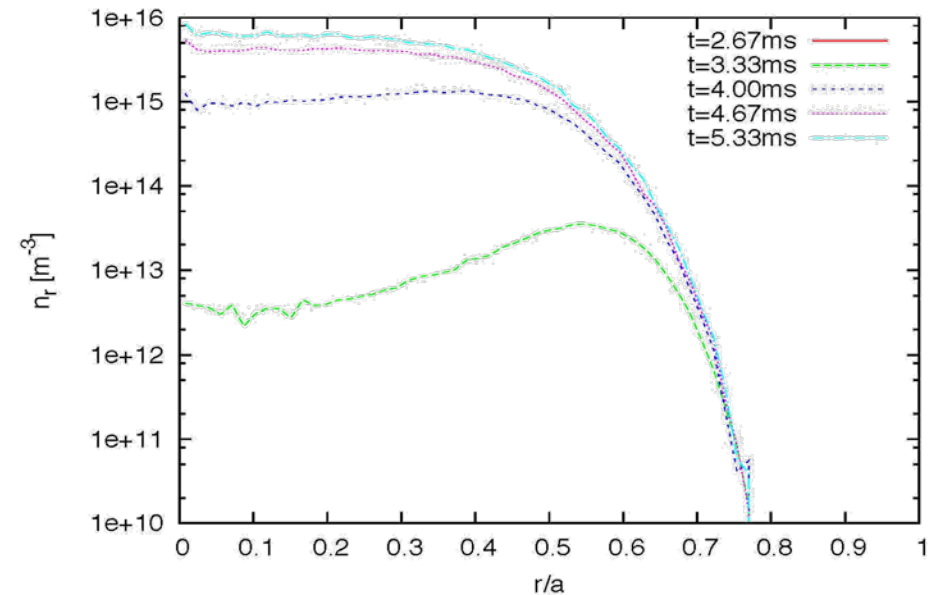
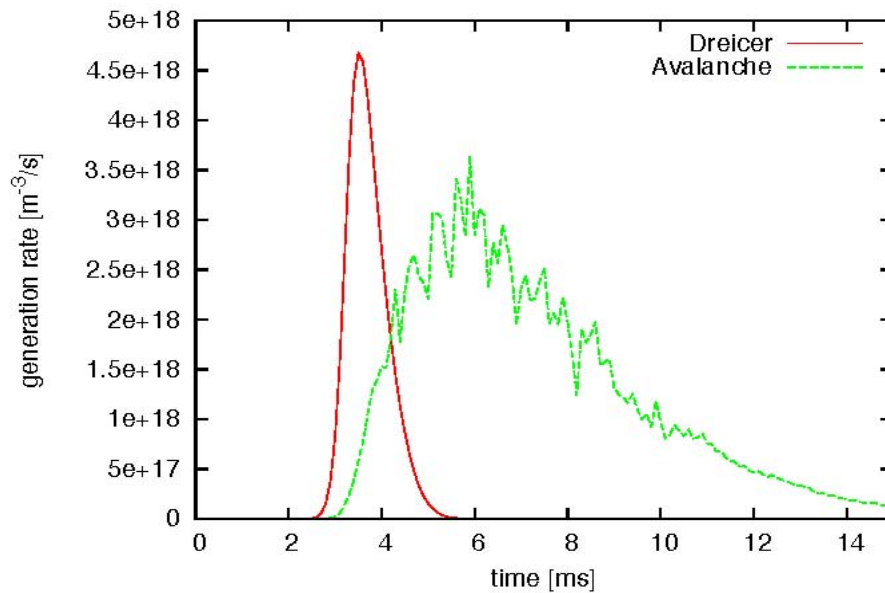
Time history of Runaway Current

- First simulation has been carried out for a JT-60U size plasma
 - elliptic cross section
 - $R = 3.4$ m, $B = 3$ T, $a = 1$ m, $\kappa = 1.6$, $I_p = 2.5$ MA
 - assuming thermal quench such that $T_e = 2$ keV $\rightarrow$ 10 eV with constant electron density of 3×10^{19} m $^{-3}$.



- Current quench ~ 10 ms
- Runaway current ~ 750 kA
- Agree well with prediction by a cylindrical modeling (Smith, et al., PoP, 2006).

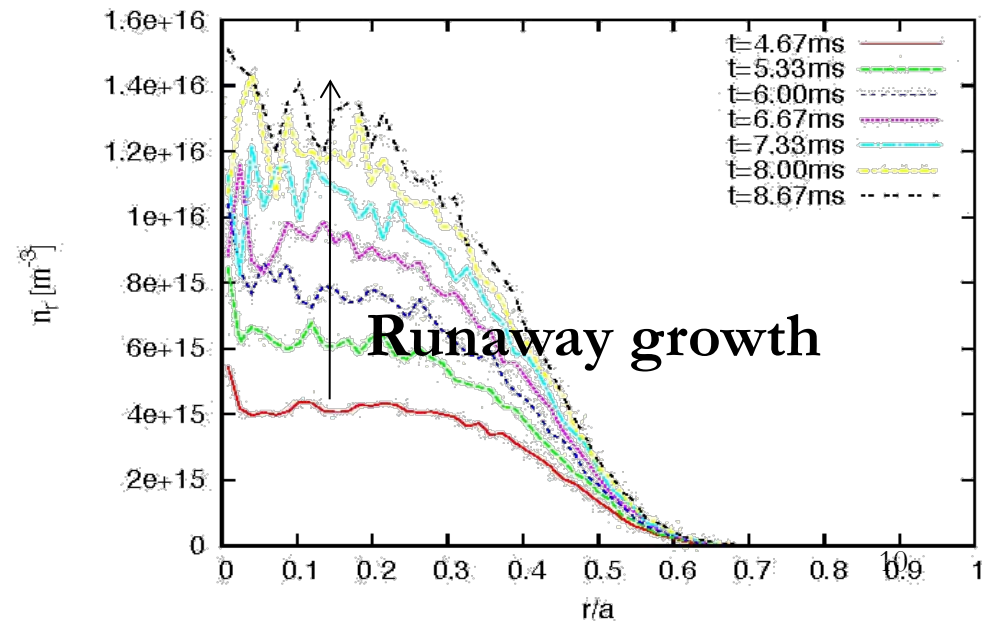
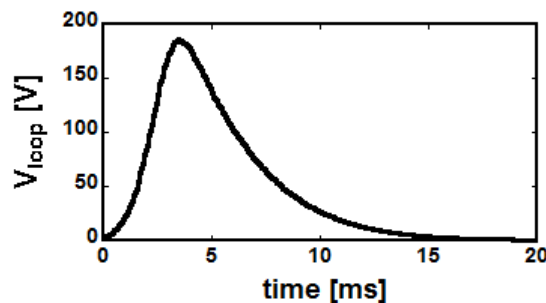
Runaway Density Profiles



- Dreicer generation is significant only in early phase of disruption.

$T < 5ms$: Dreicer generation

$T > 5ms$: Avalanche

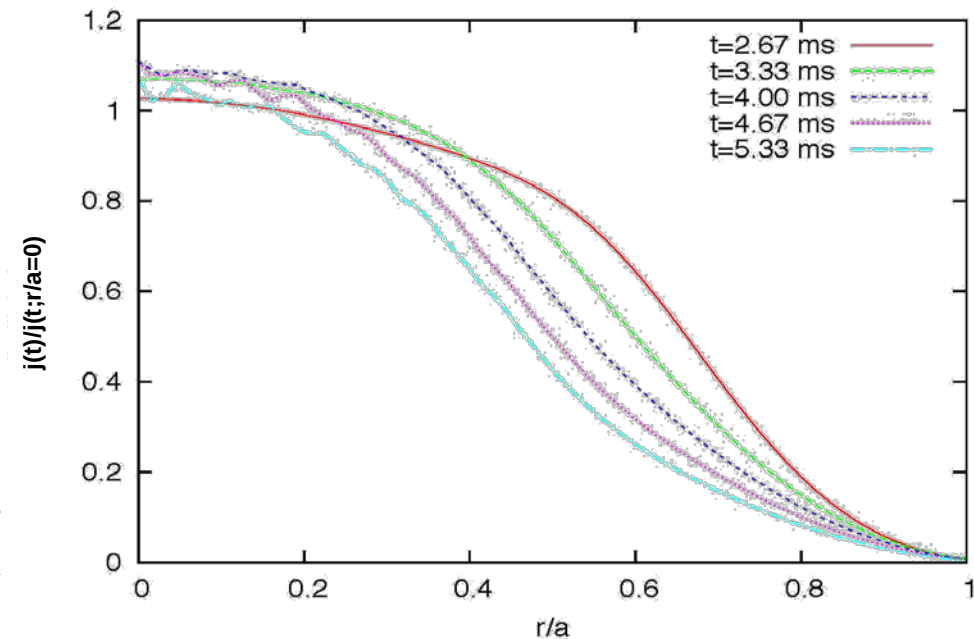
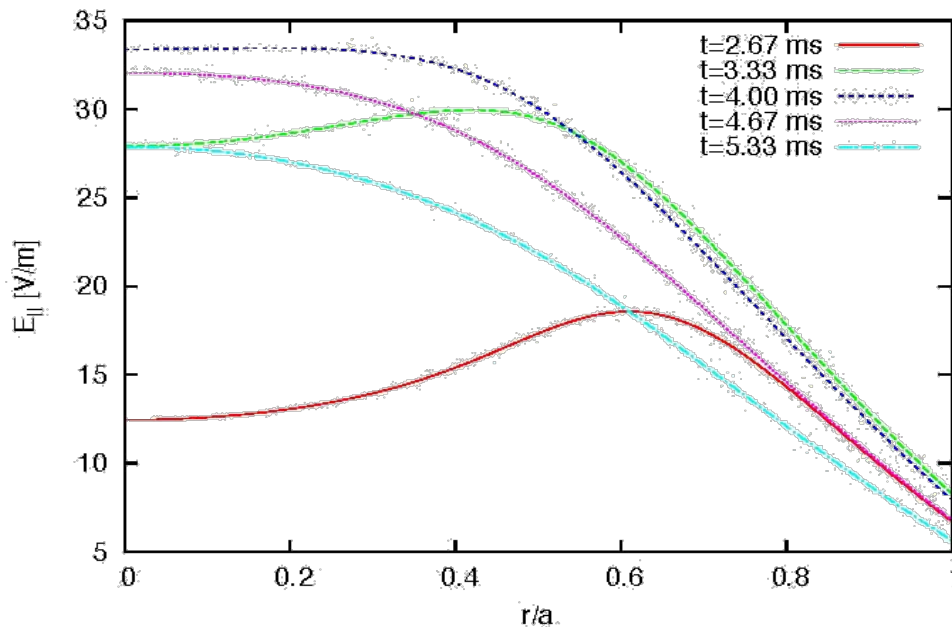


Current Profile Peaking with Runaways

- Electric field and current profiles

- Current profile peaking is consistent with Eriksson PRL 2004.
- Electric field threshold: $E_c \sim 10^{-2}$ V/m, $E_d \sim 10^3$ V/m.

Only hot tail of electrons becomes runaways: $E_c \ll E_{||} \ll E_d$



Summary

- A new version of ETC-Rel has been developed, including runaway generation mechanisms and self-consistent electric field.
- In present version, **instead of directly solving the Dreicer acceleration, generation processes are included by analytical expressions of the runaway source** to save the computational time. First simulation results for JT-60U grade disruption have been presented.
- In future, **effects of low-order magnetic perturbation (such as $n = 1$) will be investigated, which cannot treat by conventional bounce-averaging approach because of their 3-D nature.**
- Integration with 3-D MHD codes is currently being planned.