# Multi-Time-Scale Energetic Particle Dynamics in JT-60U Simulated with MHD Activity, Sources and Collisions

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Discussions with K. Tani (NAT, Japan) are thankfully acknowledged.

This work is partially supported by JSPS Grants.

The simulations were carried out using the supercomputer system HELIOS at IFERC, Aomori, Japan, under the BA collaboration between Euratom and Japan, implemented by F4E and JAEA.







20th NEXT Workshop, Kyoto, Japan, January 13-14, 2015

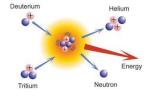
#### PREAMBLE:

**Nuclear Fusion Research at JAEA** 

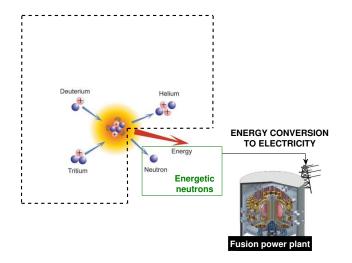
and

20 Years of Numerical EXperimental Tokamak (NEXT) Project

# Fusion reaction most easily realized in lab: $T(d,n)^4He$

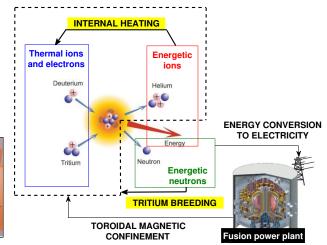


# Final goal: Electric power plants driven by D-T fusion



Figures: iter.org (reaction, power plant)

# Mainstream: Self-sustained "burning" tokamak plasma

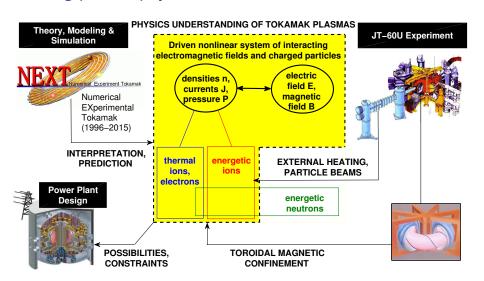


Tokamak:



Figures: iter.org (reaction, power plant)
"Harnessing the Power of Stars" J.Tachon, P.-J. Paris (tomamak)

# Burning plasma physics research at JAEA



Figures: www-jt60.naka.jaea.go.jp (JT-60U)
"Harnessing the Power of Stars" J.Tachon, P.-J. Paris (tomamak)

# Numerical EXperiment of Tokamak (NEXT) Project

NEXT Numerical Exceriment Tokamak

Started in 1996 at JAEA, the main objectives of the NEXT project\* are

- 1. to understand complex physical processes in present-days and next-generation tokamak plasmas,
- to predict and evaluate the plasma performance of tokamak reactors, such as ITER (International Thermonuclear Experimental Reactor), and
- 3. to contribute to the progress in plasma physics and related research areas via numerical simulation.

#### Focus of this talk:

Advances in the study of interaction between MHD waves and fast ions.

<sup>(\*)</sup> http://www-jt60.naka.jaea.go.jp/english/theory/intro/index.html

#### **INTRODUCTION:**

# MHD Waves and Energetic Ions in Tokamak Plasmas

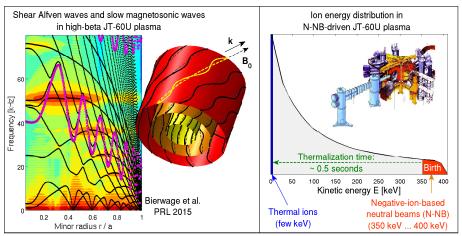
#### JT-60U plasma in reactor-relevant parameter regime

#### High plasma beta:

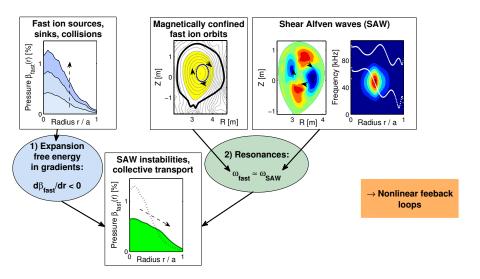
$$\frac{2\mu P}{\beta = \frac{2\mu P}{B^2} = \frac{\text{thermal pressure}}{\text{magnetic pressure}} \approx 3.5\% \qquad \text{(power plant: } \beta \sim 5\%\text{)}$$

• Fast ions generated by powerful energetic particle beams:

up to 400 keV, 5 MW,  $\beta_{\rm fast} \sim \beta_{\rm total}/2$  (power plant: 3.5 MeV <sup>4</sup>He)



# Physical effects governing fast ion dynamics in core plasma



⇒ Challenges: multiple time scales and complicated nonlin. interactions

#### Current research activities

- Study dynamics of fast-ion-driven modes.
- Analysis of linear and nonlinear resonant wave-particle interactions.
- Fast ion transport and bulk heating.
- Comparison between simulations and experiments (V&V).

#### Current research activities discussed in this talk

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"Multi-Time-Scale Energetic Particle Dynamics

Dynamics of Energetic Particle Modes
in JT-60U Simulated with MHD Activity,

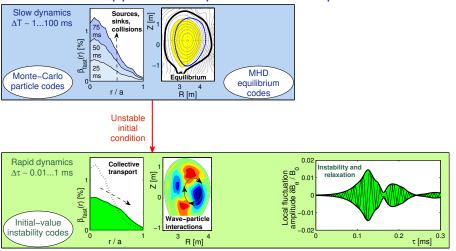
Sources and Collisions"

#### Outline

"Multi-Time-Scale Dynamics of Energetic Particle Modes in JT-60U"

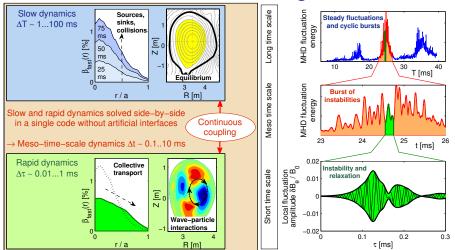
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# Conventional approach: Separation of temporal scales



- Pros: Clear separation of physical effects.
   Trends revealed by varying free parameters.
   [e.g., Bierwage et al., Nucl.
   Fusion 54 (2014) 104001]
- <u>Cons:</u> Overestimation of fast ion confinement in MC codes.
   Overestimation of fluctuation level and transport in instability codes<sub>10/25+5</sub>

#### Advanced method: Self-consistent long-time simulation



- **Pros:** Meso-*t*-scale dynamics covered. Bridge between models and experiments.
- **Cons:** Complicated. Computationally expensive.

[2014 IAEA FEC: Todo TH/7-1, Bierwage TH/P7-39]

# Hybrid model for MHD and fast ion dynamics in MEGA

• Bulk plasma modeled as 3-D MHD fluid: full MHD model,  $(R, \varphi, Z)$  finite difference grid, 4th-order Runge-Kutta solver

$$\rho_{\rm m} \partial_t \mathbf{V} = -\rho_{\rm m} \left[ \mathbf{\Omega} \times \mathbf{V} + \frac{1}{2} \nabla V^2 \right] + (\mathbf{J} - \mathbf{J}_{\rm fast}^{\rm eff}) \times \mathbf{B} - \nabla P \qquad (1)$$
$$-\rho_{\rm m} \left[ \nu \nabla \times \mathbf{\Omega} - \frac{4}{3} \nu \nabla (\nabla \cdot \mathbf{V}) \right],$$

(2)

(3)

(4)

$$egin{aligned} \partial_t 
ho_{\mathrm{m}} = & - oldsymbol{
abla} \cdot (
ho_{\mathrm{m}} oldsymbol{
abla}), \ \partial_t P = & - oldsymbol{
abla} \cdot (P oldsymbol{
abla}) - (\Gamma - 1) P 
abla \cdot oldsymbol{
abla} 
\end{aligned}$$

 $\partial_{t} \mathbf{B} = -\nabla \times \mathbf{E}$ .

$$+\left(\Gamma-1
ight)\left\{\eta J^2+
u
ho_{
m m}[\Omega^2+rac{4}{3}(oldsymbol{
abla}\cdotoldsymbol{{\sf V}})^2]
ight\};$$

with  $\mathbf{E} = -\mathbf{v} \times \mathbf{B} + \eta \mathbf{J}$ ,  $\mu_0 \mathbf{J} = \mathbf{\nabla} \times \mathbf{B}$ ,  $\Omega = \mathbf{\nabla} \times \mathbf{v}$  and no-slip boundary.

• Fast ions modeled as guiding centers: PIC method,  $\mathrm{d}\mu/\mathrm{d}t=0$ 

$$d_t \mathbf{R}_{gc} = \mathbf{v}_{\parallel}^* + \mathbf{v}_{E} + \mathbf{v}_{B}, \qquad m v_{\parallel} d_t v_{\parallel} = \mathbf{v}_{\parallel}^* \cdot [e_{H} \mathbf{E} - \mu \nabla B] \qquad (5)$$

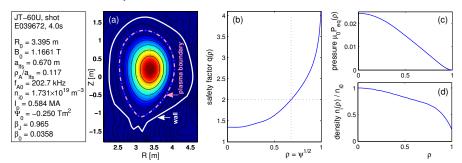
with  $\mathbf{v}_{\parallel}^* = v_{\parallel} [\mathbf{B} + \rho_{\parallel} B \nabla \times \mathbf{b}] / B^*$ ,  $\rho_{\parallel} = v_{\parallel} / \Omega_{\mathrm{L}}$ ,  $B^* = B(1 + \rho_{\parallel} \mathbf{b} \cdot \nabla \times \mathbf{b})$ .

Loss boundary condition at wall. FLR simulated via satellite particles.

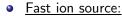
Trubnikov collision operators (slowing down, scattering, diffusion).

MEGA: Todo et al, Phys. Plasmas 12 (2005) 012503, Phys. Plasmas 13 (2006) 082503<sub>5+5</sub>

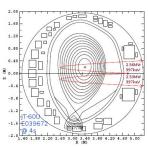
# Simulation setup: JT-60U shot E039672 @ t = 4s



- MHD equilibrium and profiles:
  - ▶ B(R, Z), P(R, Z) reconstructed numerically
  - enforce  $\nabla P_{\text{eq}} = \mathbf{J}_{\text{eq}} \times \mathbf{B}_{\text{eq}}$ , fit  $n_{\text{i}}^{\text{exp}}$
- Coupling between MHD and fast ions:
  - include only n = 1 harmonic (wavelength = torus circumference) to focus on n = 1 modes observed in exp.



deposition profile computed for pair of tangential N-NBs



#### Outline

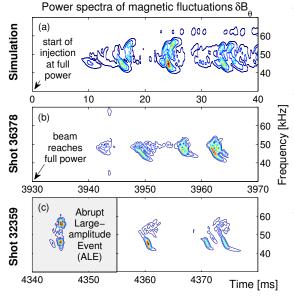
"Multi-Time-Scale Dynamics of Energetic Particle Modes in JT-60U"

1. Simulation model and methods:

Advances towards predictive simulations for energetic particle dynamics

- Meso- and long time scale (0.5 ms 40 ms): Reproduced chirping bursts of JT-60U exp.
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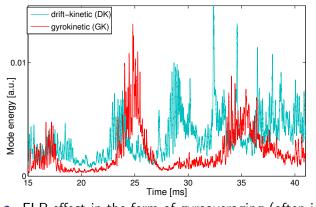
# MHD dynamics on 0.5 ms - 40 ms scale V&V: Reproduced chirping bursts of JT-60U

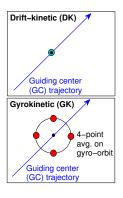


- Robust features:
  - ► down-/up-chirping in 40-60 kHz band.
    - few ms life time,
  - ▶ periods of 5-15 ms.
- Agreement shows that:
  - essential mechanisms of fast ion dynamics and plasma response are modelled correctly,
  - sim. results are relevant for rapidly chirping bursts.
- Note:
  - sim. measures signals at the mode location
  - exp. measurements are done outside of the plasma

15/25 + 5

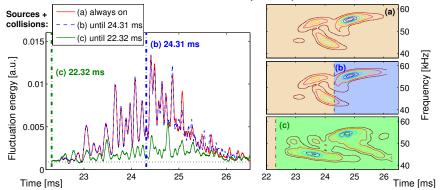
# FLR effect (gyroaveraging) proved to be essential for reproducing intermittency of bursts





- FLR effect in the form of gyroaveraging (often ignored previously) reduces wave-particle coupling strength and leads to
  - ► Longer quiet periods between bursts.
  - ► Lower fluctuation levels during quiet periods.
- ⇒ New mechanism (besides damping) that controls intermittency! [For damping effect, see Todo *et al.*, *Nucl. Fusion* **52** (2012) 033003]

# Effect of sources and collisions (S&C) on bursts



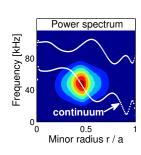
- Results of turning S&C off at (b) peak or (c) start of a burst:
  - ▶ (b): Decay phase and chirping indep. of S&C.
  - ► (c): S&C is essential for rise of base level of fluctuation energy, but even without S&C burst and chirps still occur at lower amplitude.
- Implications:
  - ▶ No retardation. ⇔ Dominant resonances are near birth energy.\*
  - ▶ We may learn sth. about deviation from marginal stability. ⇒ Study more!

<sup>(\*)</sup> Bierwage & Shinohara, *Phys. Plasmas* **21** (2014) 112116

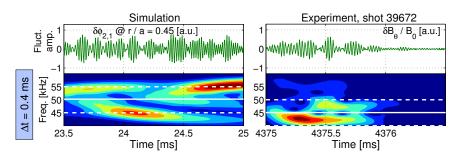
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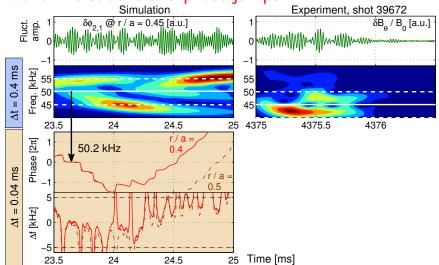
#### MHD dynamics on 10 $\mu$ s - few ms scale



- Previously reported chirping dynamics based on 0.4-1 ms FT windows.
- However: Time traces of raw signal exhibit pulsations on shorter scale.

#### MHD dynamics on 10 $\mu s$ - few ms scale

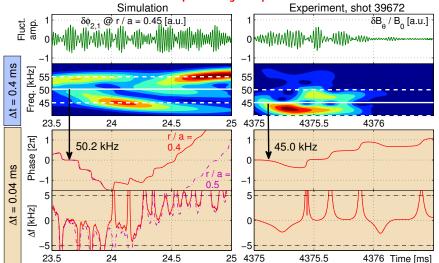
Found: Pulsations and phase jumps



• Analysis on 0.04 ms scale reveals: Phase jumps  $(\pm \pi)$  between pulses. Wave freq. varies less  $(\pm 2.5 \text{ kHz})$  than averaged chirps  $(\pm 5 \text{ kHz})$ .

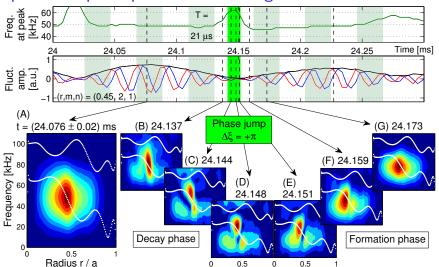
#### MHD dynamics on 10 $\mu$ s - few ms scale

Found: Pulsations and phase jumps



- Pulses and phase jumps discovered in sim. are also found in exp. data.
- Not considered in theories for chirping.  $\Rightarrow$  Improve physical picture!<sub>19/25+5</sub>

# Spatio-temporal pulsation of Energetic Particle Modes\*



- Resonant drive at certain freq. keeps phases of dispersive waves aligned.
- Complicated interaction between mode structure and fast ion distrib.

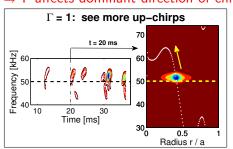
# Struct. of Shear Alfvén Wave continuum affects chirping

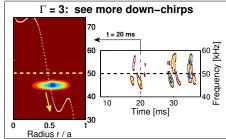
For ideal monoatomic gas (3 degrees of freedom, isotropic collisions)

 $\Gamma = C_p/C_V = 1 + 2/F = 5/3$ . For strongly magnetized tokamak plasma (electrons, deuterons, high-Z impurities) appropriate  $\Gamma$  value not known.

$$\partial_t P = -\nabla \cdot (P\mathbf{V}) \underbrace{-(\mathbf{\Gamma} - 1)P\nabla \cdot \mathbf{V}}_{\text{compressional response}} \dots \Rightarrow \underbrace{\text{controls continuous}}_{\text{SAW spectrum}}$$

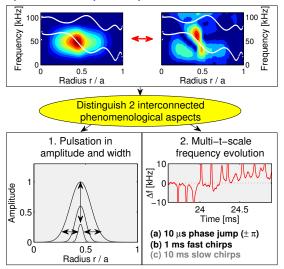
 $\Rightarrow \Gamma$  affects dominant direction of chirping:



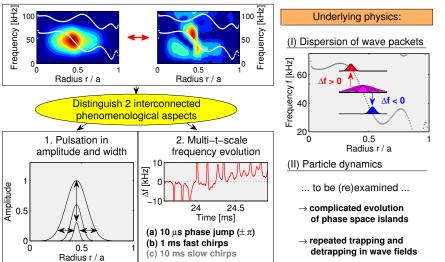


- Dominant component of signal tends to chirp towards nearest accumulation point, where  $d\omega_A(r)/dr$  is small.
- One out of several new ingredients for a complete theory of chirping. 21/25+5

#### Spatial-temporal pulsation of EPM: Summary



# Spatial-temporal pulsation of EPM: Summary



#### • To complete the picture: Clarify self-consistent evolution of drive.

Revisit theory of phase space dynamics [Berk et al., Phys. Lett. A 234 (1997) 213], which assumed small deviation from marginal stability and dominance of discrete eigenmodes with constant mode structure, ignoring continuous spectra.

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#### Summary and conclusion

#### • Major numerical advances:

 Developed self-consistent simulation framework with MHD, sources and collisions: Promises new insights and better predictions!

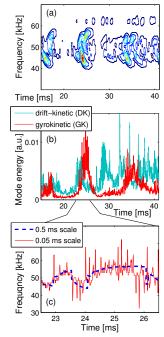
#### • Multi-time-scale simulation results:

Reproduced rapid chirps (~ 1 ms) from experiments and demonstrated importance of fast ion FLR effects:
 Successful V&V!

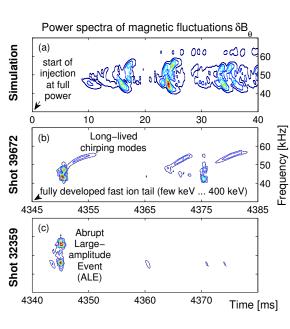
3. Analysis of chirps revealed spatio-temporal pulsations of EPMs ( $\sim 100\,\mu s$ ) and phase jumps ( $\sim 10\,\mu s$ ) in both sim. & exp.: Sucessful V&V. New building blocks for more complete theory of chirping modes.

#### Open questions:

- Reason for  $\pm \pi$  phase jumps?
- ► Particle dynamics during EPM pulsation?



#### Outlook



- Physics:
  - Wave-particle interactions.
  - Burstiness and marginal stability.
  - Kinetic effects of bulk ions, bulk ion heating.
  - V&V, Prediction:
    - Apply to other machines and other modes.
- Long-lived chirping modes:
  - ► lasting about 5 ms
  - sensitive to plasma equil.may require fully
  - developed fast ion tail (below 300 keV)
- Abrupt Large Events
  - 4-10 times larger  $\delta B_{\theta}$
  - ▶ may involve n > 1 modes\*
- (\*) Bierwage *et al*, *Nucl. Fusion* **54** (2014) 104001

25/25+5

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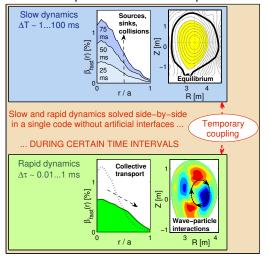
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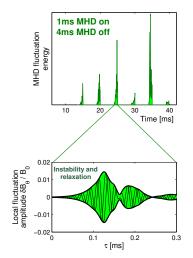
# A1. Numerical parameters and computational resources

- Setup discussed here:
  - ► MHD fluctuations with single toroidal mode number *n* = 1 interacting with 2 fast ion beams
  - Step size:  $\Delta t_{\rm MHD} = 0.05 \, \tau_{\rm A0} = 0.04 \, \mu \text{s} \, (\Delta t_{\rm PIC} = 4 \times \Delta t_{\rm MHD})$
  - Spatial grid:  $N_R \times N_{\omega} \times N_Z = 384 \times 32 \times 352$
  - ► MPI domains:  $M_R \times M_{\varphi} \times M_Z = 32 \times 4 \times 32$  (4096 MPI processes on 256 nodes with 16 cores each)
- Simulation of 50 ms physical time:
  - ▶  $13.9 \times 10^6$  time steps, accumulating  $3.38 \times 10^6$  simulation particles
  - ▶ Wall time: 11.3 days for 25 ms, 26 days for 50 ms (160k node hours)



#### A2. Compromise: Multi-phase simulation method

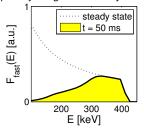




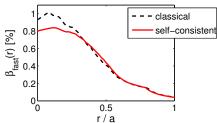
- Pros: Efficient and useful for prediction of realistic fast ion profiles.
  - [Todo et al., Nucl. Fusion **54** (2014) 104012]
- Cons: Only rudimentary meso-t-scale dynamics ( $\lesssim 1$  ms).

# A3. Effect of MHD activity on fast ion distribution

Spatially integrated velocity distribution



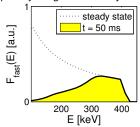
Velocity-integrated radial profile t = 50 ms

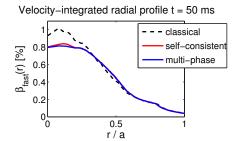


- Classical (no MHD): Overestimates fast ion gradients ⇒ unstable.
- Self-consistent: Fast ion energy in central core reduced by  $\approx 15\%$ .

# A3. Effect of MHD activity on fast ion distribution

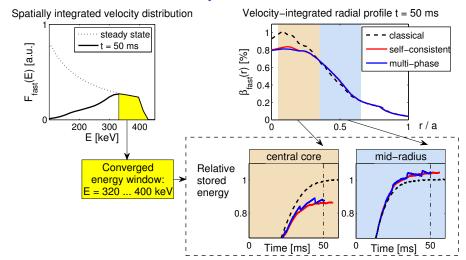
Spatially integrated velocity distribution





- Classical (no MHD): Overestimates fast ion gradients  $\Rightarrow$  unstable.
- Self-consistent: Fast ion energy in central core reduced by  $\approx 15\%$ .
- Multi-phase (MHD 1ms on/4ms off): Reproduces self-consist. sim. result.

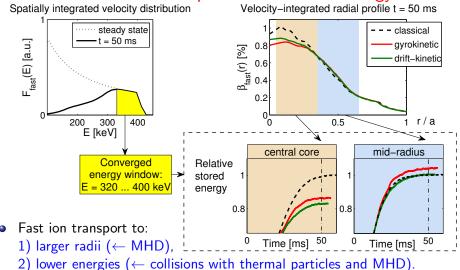
# A3. Effect of MHD activity on fast ion distribution



- Classical (no MHD): Overestimates fast ion gradients ⇒ unstable.
- Self-consistent: Fast ion energy in central core reduced by  $\approx 15\%$ .
- Multi-phase (MHD 1ms on/4ms off): Reproduces self-consist. sim. result.

# A4. Effect of MHD actitivty on fast ion distribution

# Found: FLR affects transport in radius and energy



- Drift-kinetic (no FLR): Underestimates radial transport by \$5%.
  - Overestimates energy transfer to bulk plasma by  $\lesssim 5\%$ . For n = 1! 30/25+5