



# Gyrokinetic analysis of turbulent transport in helical systems with different magnetic shear

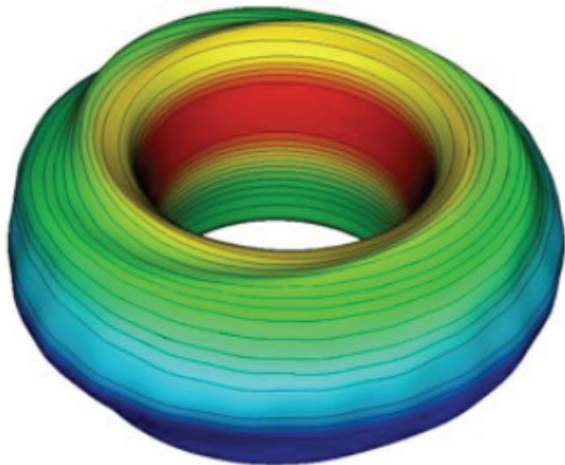
A. Ishizawa  
Kyoto University

# Motivation

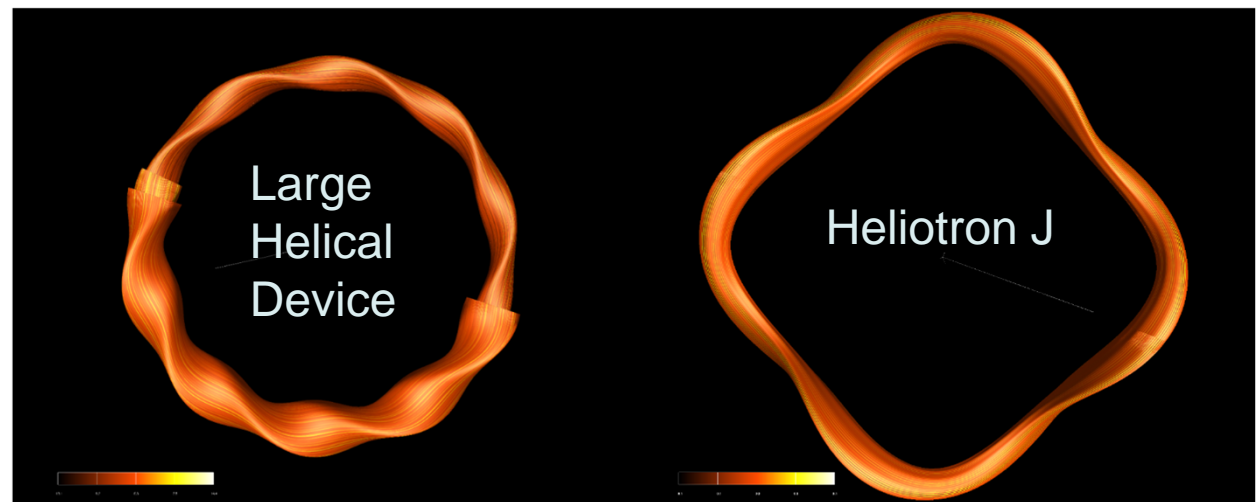
- Approximate degrees of deviation from axisymmetry
  - **3D tokamak** ( $\delta B_{n \neq 0} / \delta B_{n=0} = 10^{-3}$  to  $10^{-2}$ )
  - **Helical reversed field pinch state** ( $\delta B_{n \neq 0} / \delta B_{n=0} = 0.03$  to  $0.05$ )
  - **Stellarator** ( $\delta B_{n \neq 0} / \delta B_{n=0} = 0.1$  to  $0.3$ )

D. Spong, APS 2014

Tokamak with RMP  
Resonant fields - ELM coil



D. Spong, APS 2014



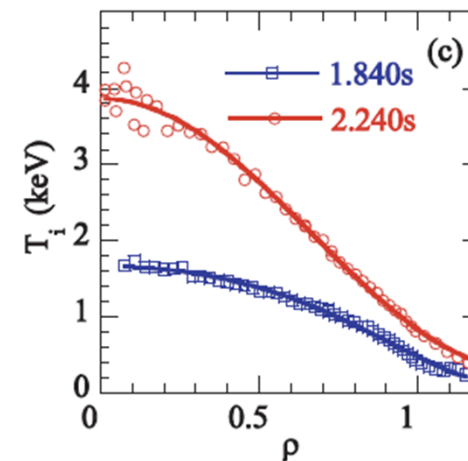
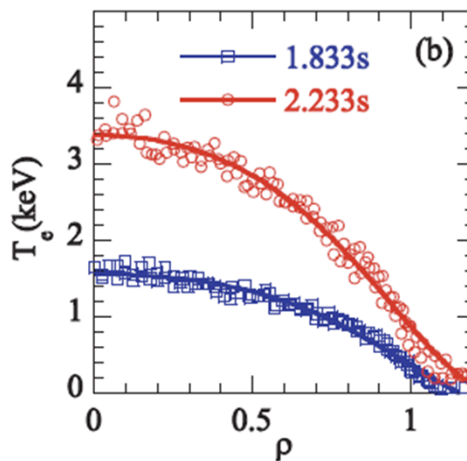
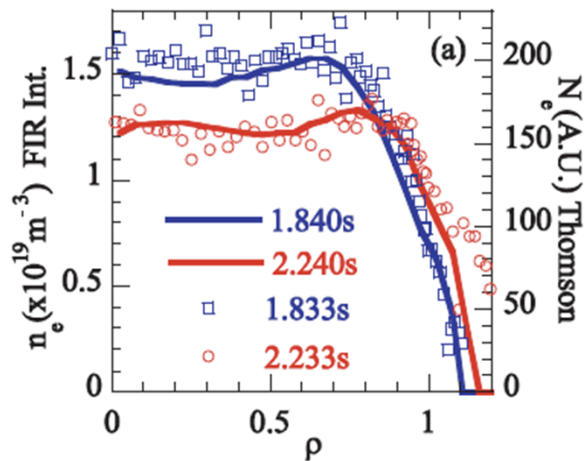
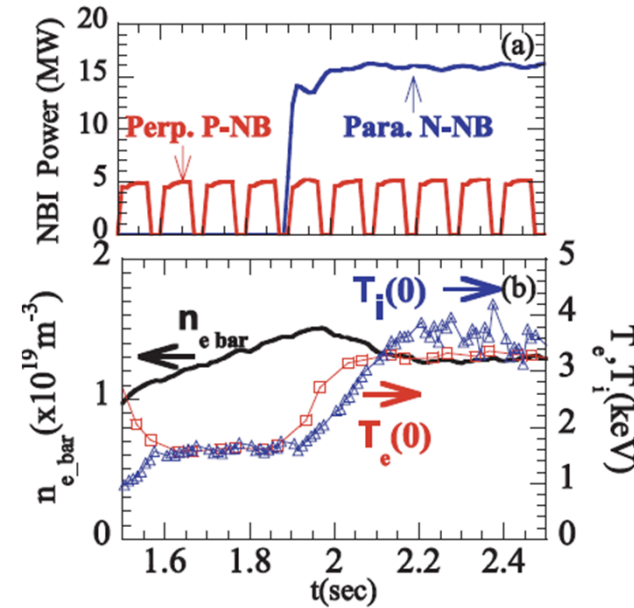
# Outline

- Introduction
- Large Helical Device (LHD)
- Heliotron J (HJ)
- Comparison
- Summary

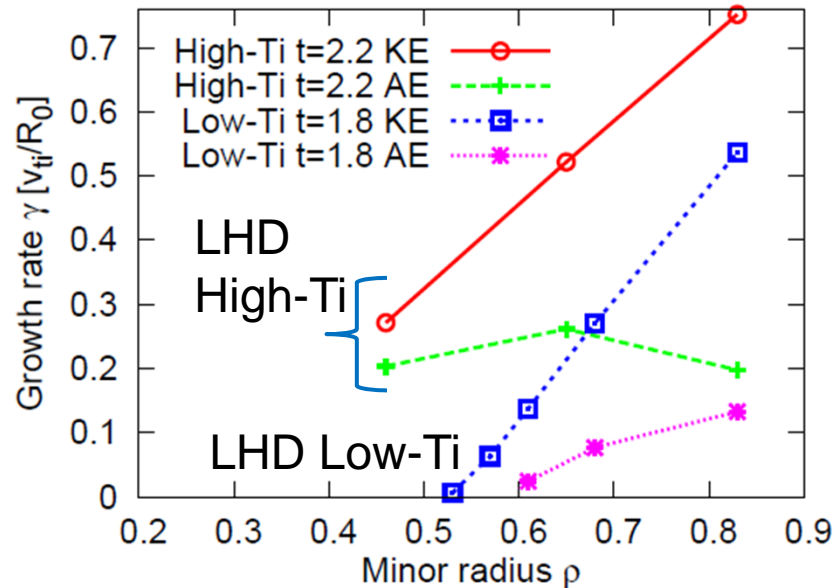
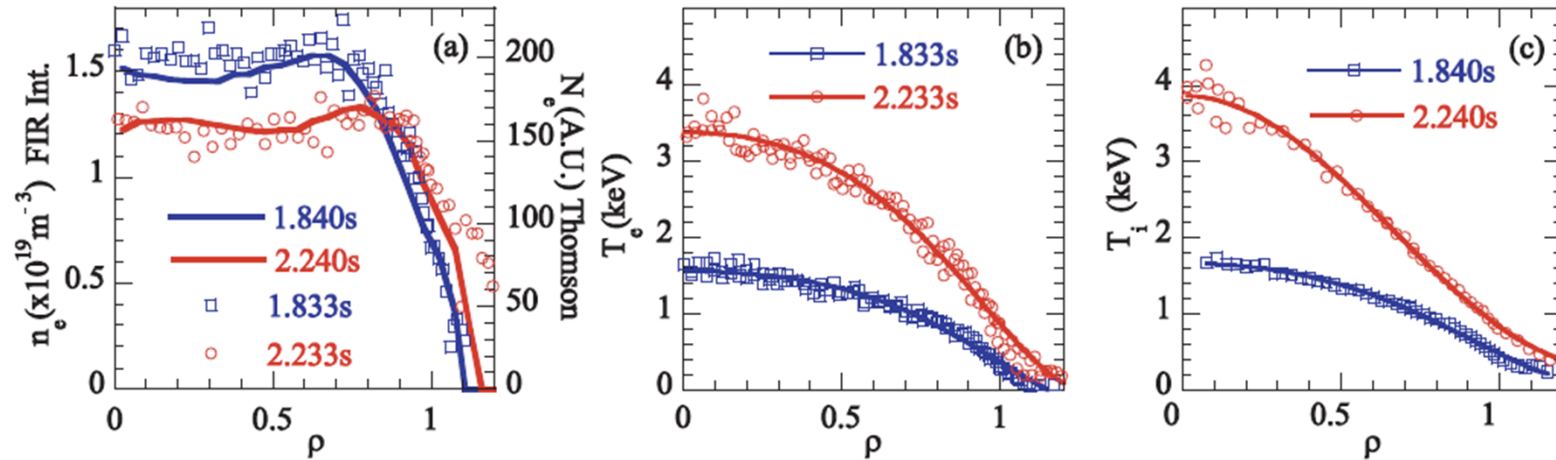
	LHD-L	HJ-ST
$R_0/a$	6.2	7.3
$\rho = r/a$	0.68	0.5
$q$	1.5	1.7
$\rho_* [10^{-3}]$	2.	4.5
$v_i^*$	0.083	3.2
$\beta [\%]$	0.2	0.05
$T_e/T_i$	0.96	1.3
$R_0/L_n$	2.7	9.3
$R_0/L_{Ti}$	8.7	13.
$R_0/L_{Te}$	9.1	17.
$\hat{s}$	1.2	0.023
$D_{\text{well}}$	-0.01	0.74

# LHD discharge #88343

- $B=2.75\text{T}$ ,  $R=3.6\text{m}$  (shifted to  $3.75\text{m}$ )
- Low-Ti phase:  $T_i=1.6\text{keV}$   $t=1.8\text{s}$
- High-Ti phase:  $T_i=3.9\text{keV}$   $t=2.2\text{s}$ 
  - $\text{Beta}(r/a=0.65)=0.3\%$
  - Collision:  $1/\nu u$

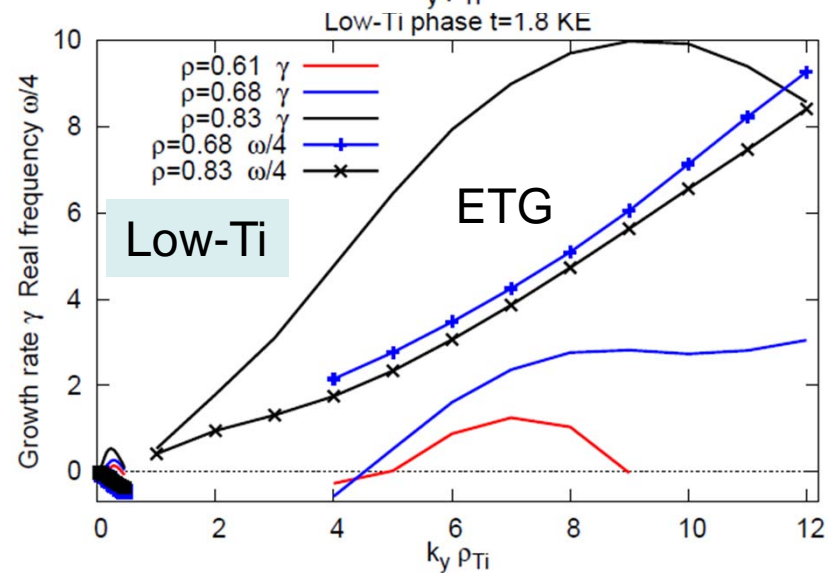
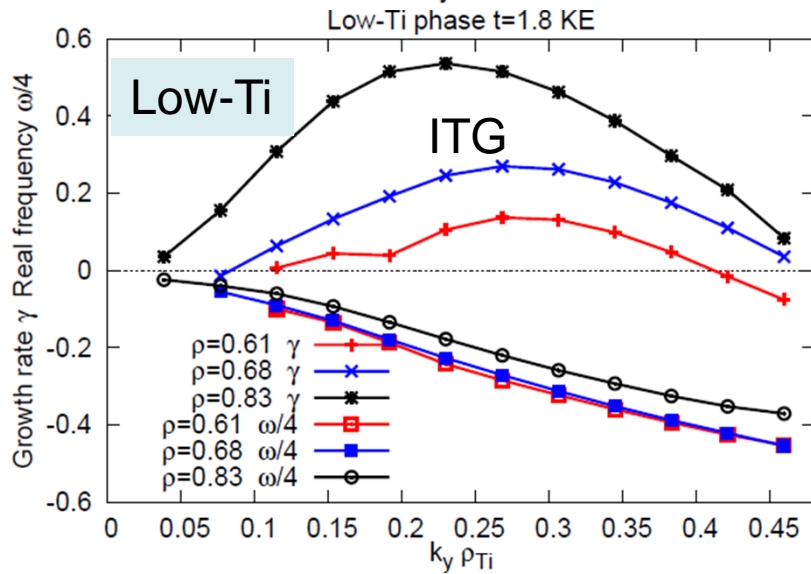
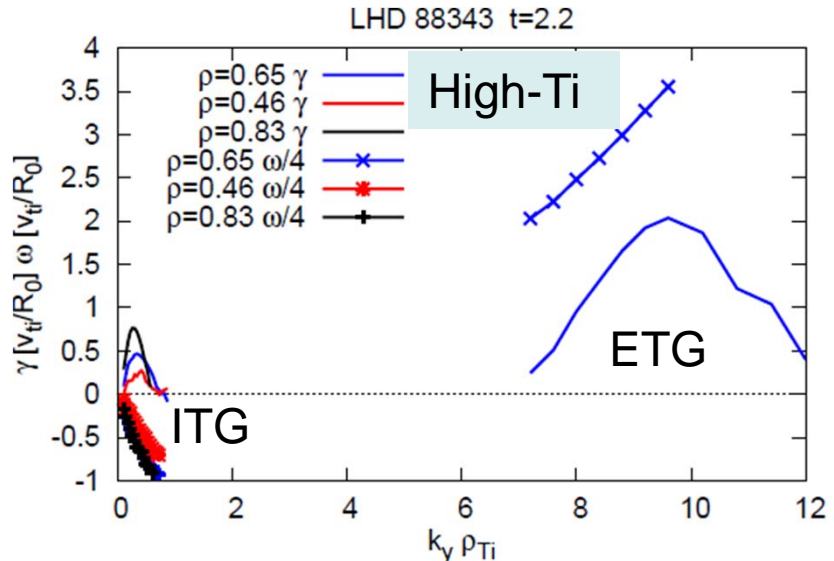
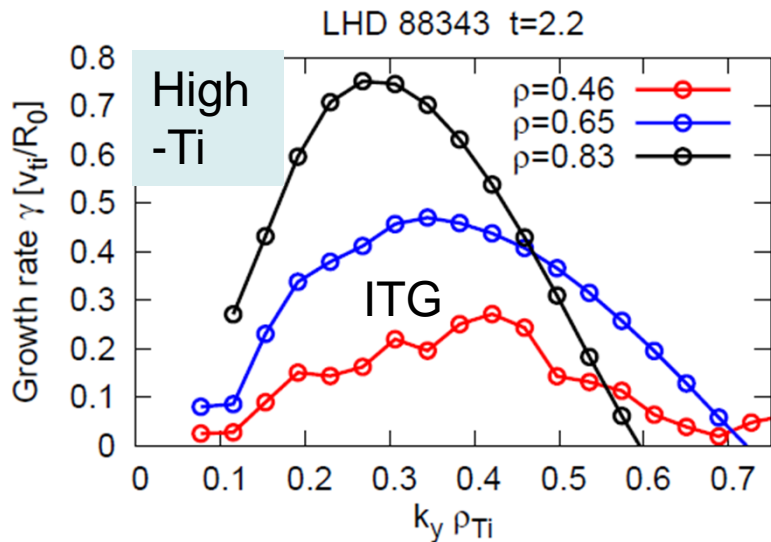


# Linear analysis of LHD



- Ion temperature gradient (ITG) modes are unstable.
- Kinetic electron effects enhance the instability.

# Linear growth rate

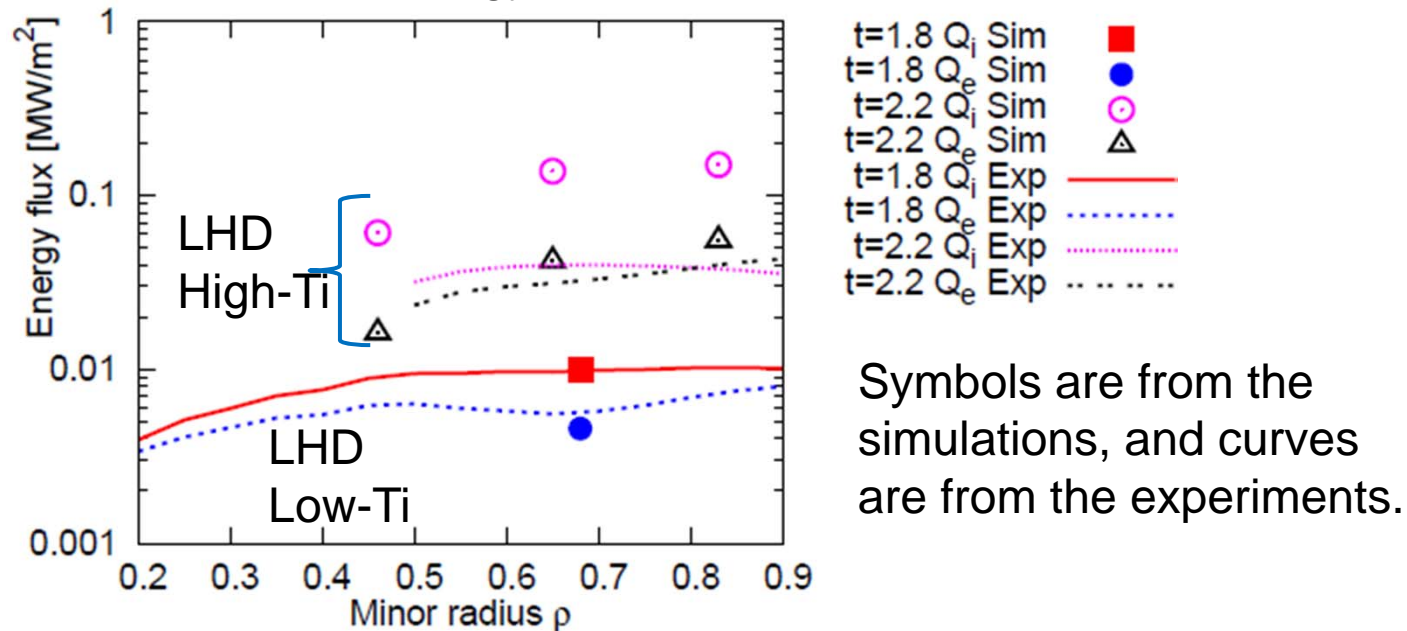


- ITG and ETG modes are unstable.



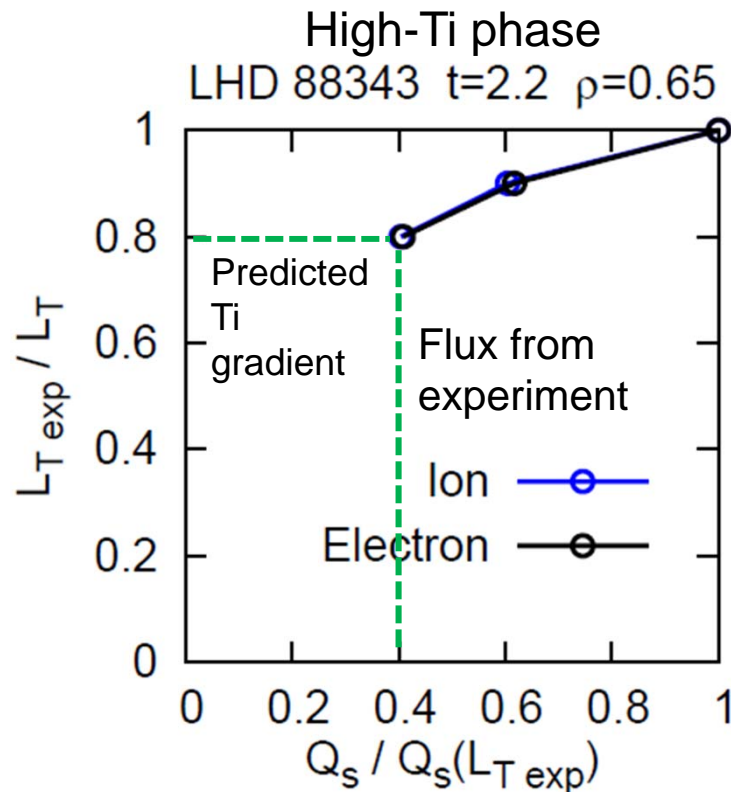
# Turbulent energy flux of LHD

Ion and electron energy fluxes due to turbulence,  $Q_i$  and  $Q_e$ .



- The transition of the energy flux from the low-Ti phase (t=1.8,  $T_i=1.6\text{keV}$ ) to the high-Ti phase (t=2.2s,  $T_i=3.9\text{keV}$ ) in the experiment is reproduced.
- There is no short-fall problem, which suffers the GK analysis of some tokamaks.

# Prediction of profile by flux matching



$L_T$  : Temperature gradient length

$L_{T \text{ exp}}$  : Experimental observation of  $L_T$

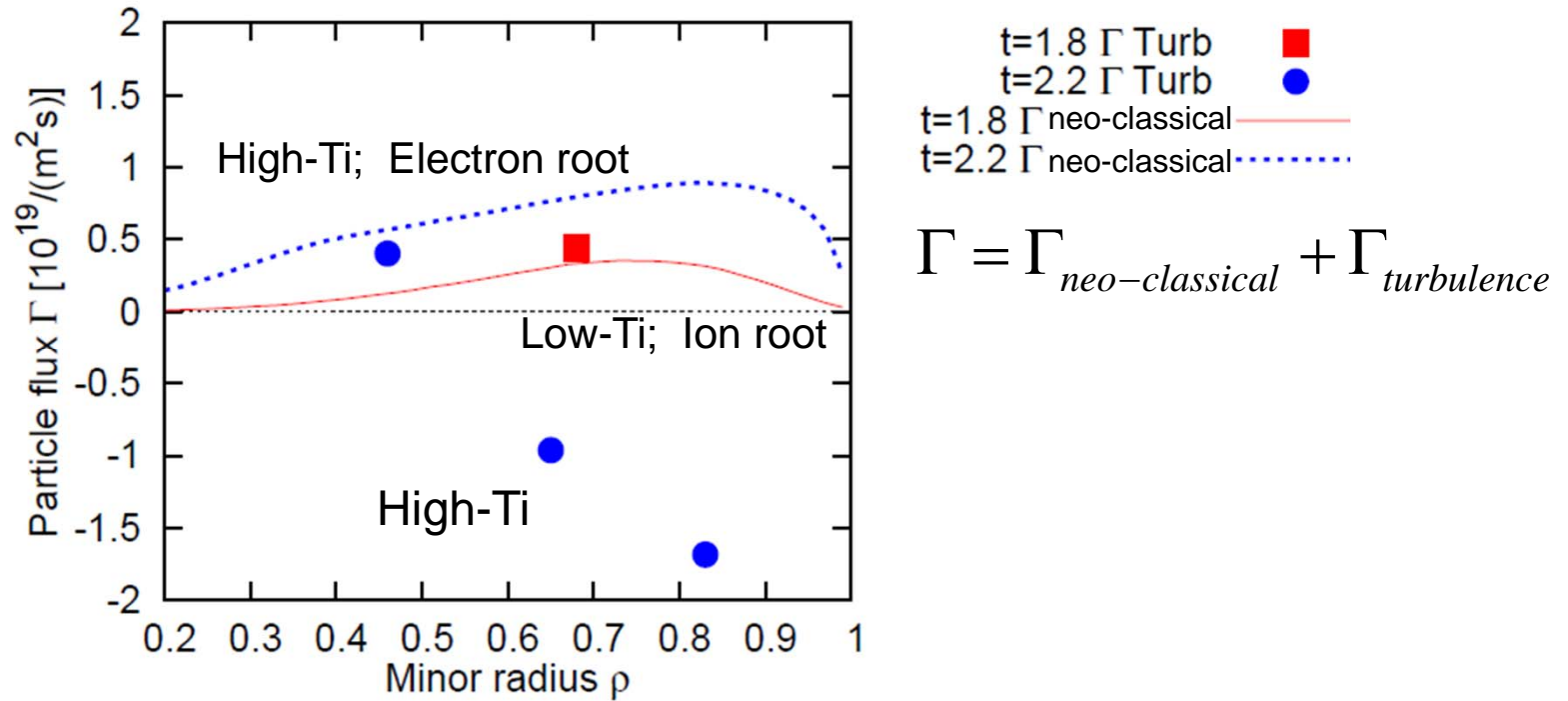
$Q_s$  : Heat flux of "s" species

$Q_s(L_{T \text{ exp}})$  : Experimental observation of  $Q_s$

- The predicted temperature gradient length deviates from experimental observation about 20%.

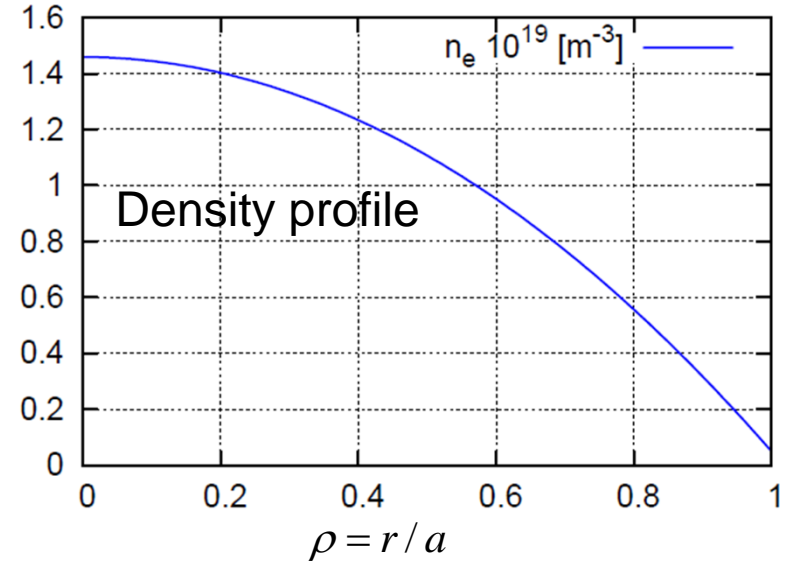
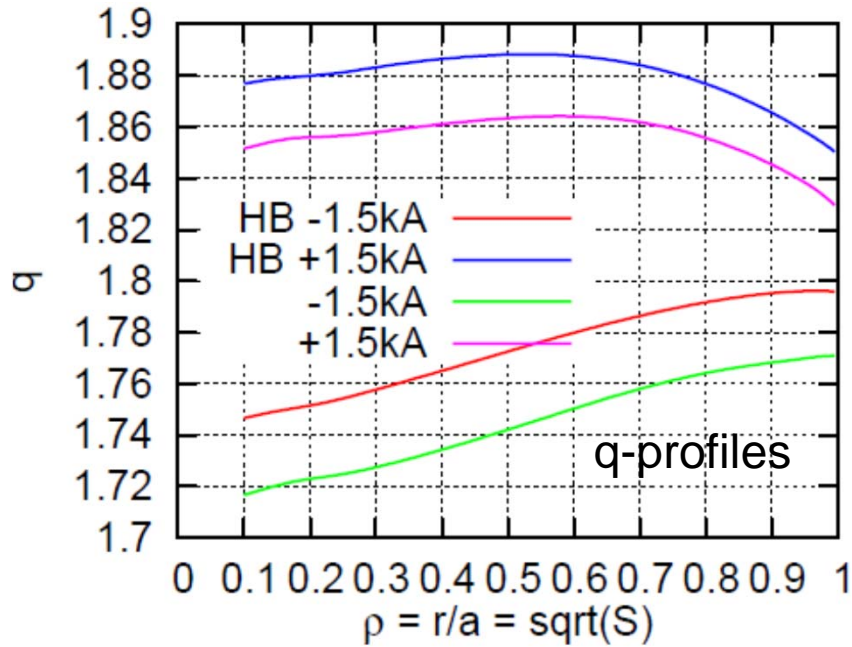


# Turbulent particle flux

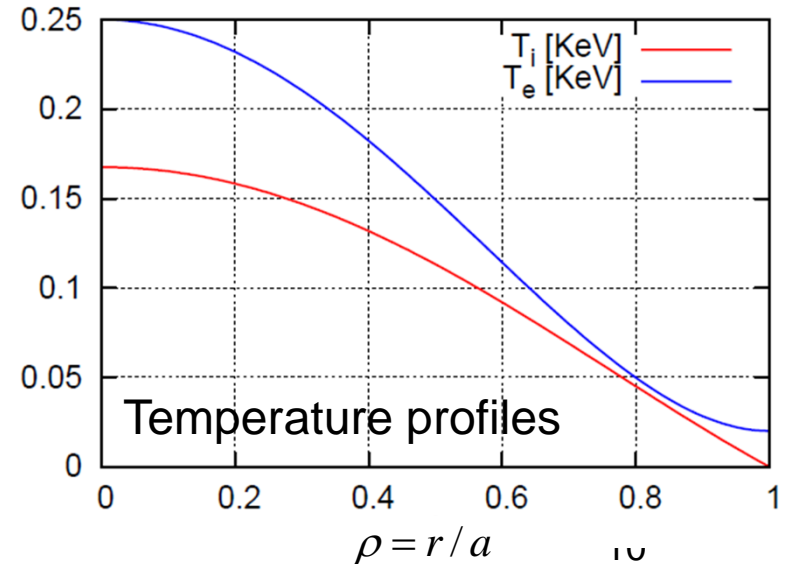
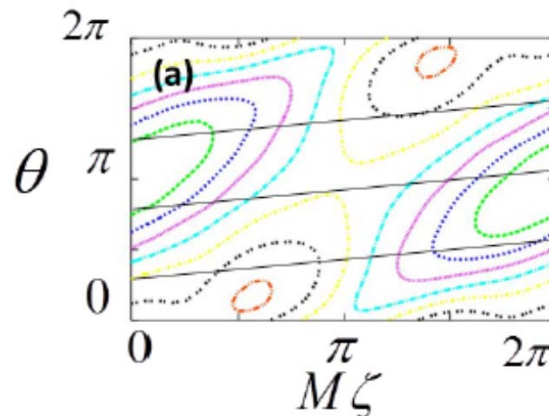
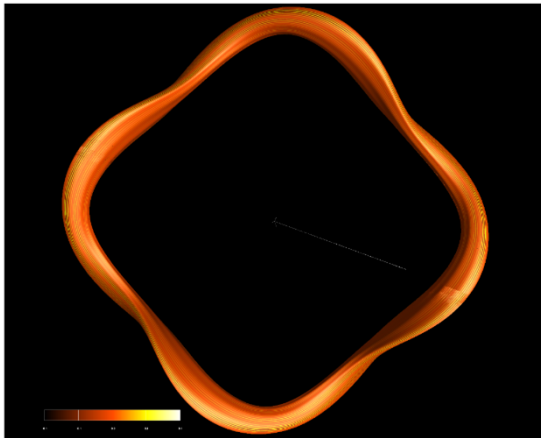


- The turbulent particle flux directs to the magnetic axis, and its direction is opposite to the neo-classical particle flux.

# Profiles of HJ plasma



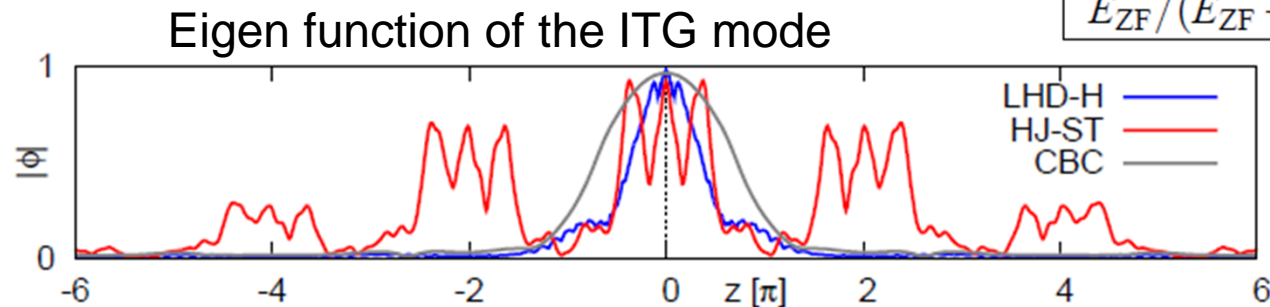
HJ-ST: standard configuration  
 HJ-HB: High-toroidal ripple configuration



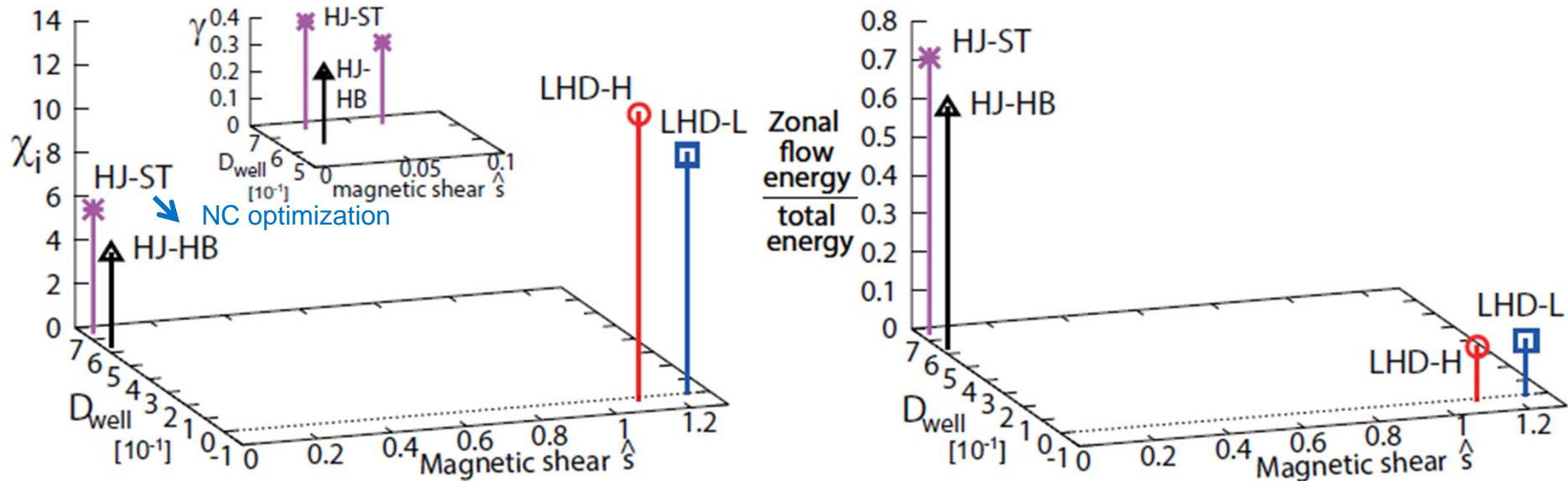
# Simulation results

- The ITG mode is unstable
  - The LHD is the inner shifted configuration, so it is magnetic hill with a moderate shear.
  - The HJ is magnetic well with a very weak shear.
  - The LHD has an advantage compared with HJ from the linear aspect of drift wave instability.

	LHD-L	HJ-ST
$R_0/a$	6.2	7.3
$\rho = r/a$	0.68	0.5
$q$	1.5	1.7
$\rho_* [10^{-3}]$	2.	4.5
$v_i^*$	0.083	3.2
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$D_{\text{well}}$	-0.01	0.74
Instability	ITG	ITG
$\gamma [v_{Ti}/R_0]$	0.27	0.4
$R_0/L_T - R_0/L_{T\text{crit}}$	2.6	5.2
$\chi_i [v_{Ti}\rho_{Ti}^2/R_0]$	11.	5.9
$\chi_e [v_{Ti}\rho_{Ti}^2/R_0]$	4.8	2.4
$E_{ZF}/(E_{ZF} + E_{ITG})$	0.14	0.72



# Comparisons in ( $\hat{s}$ , $D_{well}$ ) space



- The stabilizing effect of the magnetic shear is confirmed by the reduction of the growth rate by increasing the shear.
- The neoclassical optimization improves turbulent transport in HJ.
- Weak magnetic shear of HJ does not lead to high turbulent transport because of nonlinear interactions including zonal flow production.

# Summary

- Comparisons of turbulent transport in helical systems including the LHD and the Heliotron J
- LHD analysis
  - The simulations reproduce the temperature gradients before and after the additional NBI within 20% error.
  - There is no short-fall problem.
- HJ analysis
  - The neoclassical optimization (high-toroidal ripple) suppresses the turbulent heat transport.
- Comparison reveals that the regulation of turbulence by zonal flows is more efficient in HJ which has very weak magnetic shear.