

## Fokker-Planck simulation for runaway electron generation including the hot-tail effect

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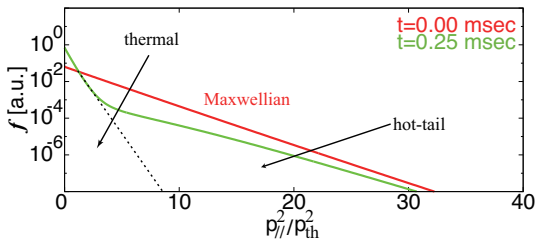


# Introduction (RE generation in disruptions)

- Runaway Electrons (REs) are generated in tokamak disruptions.
- The impact of REs to the first wall leads intolerable heat load especially in ITER.
- The high electron density plasma achieved by MGI may suppress the RE generation because of high collisionality.
- **MGI shortens the thermal quench** and might enhance the RE generation through the **hot-tail effect**.
  - This requires the kinetic treatment.
- The estimation of the amount of REs generated in tokamak disruption is required for the development of the mitigation method.

## Introduction (hot-tail effect)

- Non-thermal effect should be included for RE gen. simulation
  - If the drop of  $T$  is sufficiently fast, the plasma cools down so quickly that fast electrons do not have enough time to be thermalized.
  - The rapid cooling forms the high velocity tail of  $f_e$ .
  - It enhances **the primary RE generation** and this effect is called as **hot-tail effect**.



- For mitigated disruptions, the thermal quench time tends to be shortened.

# Equations

## Electric field diffusion

$$\frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial E}{\partial r} \right) = \mu_0 \frac{\partial j}{\partial t}$$

## Ohm's law

$$j = \sigma_{sp} E + ecn_r, \quad n_r = n_{rp} + n_{rs}$$

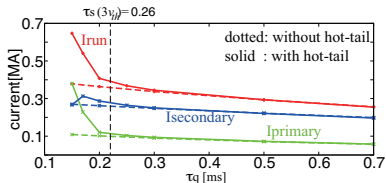
## Primary and secondary RE generation rate

$$\frac{dn_{rp}}{dt} = - \int \frac{\partial f}{\partial t} d\mathbf{p}, \quad (p_{max}^2/m \sim 0.5\text{MeV}), \quad \frac{dn_{rs}}{dt} = S_{avalanche}(n_r, E/E_C)$$

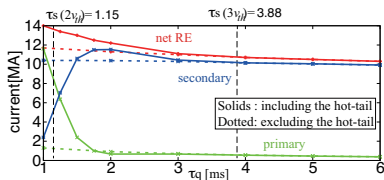
## Temperature

$$T(t) = (T(0) - T(t_{max})) \exp(-t/\tau_q) + T(t_{max}), \quad T(t_{max}) = 10\text{eV}$$

# Threshold of the hot-tail effect



JT-60U like ( $I_p=1\text{MA}$ )



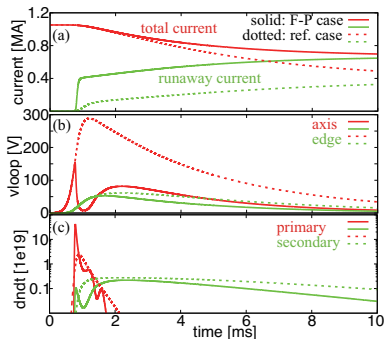
ITER like ( $I_p=15\text{MA}$ )

- If the thermal quench time  $\tau_q$  is shorter than  $\tau_s^{ee}(2 - 3v_{th})$ , the hot-tail effect becomes remarkable.

$$\tau_s^{ee}(v) = \frac{2\pi\epsilon_0^2 m_e^2 v^3}{n_e q^4 \ln \Lambda}$$

- The hot-tail effect enhances the **primary** RE current.
- The high primary RE current reduces the **secondary** RE current.

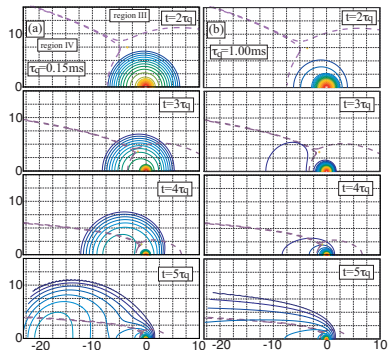
# Evolutions (JT-60U case)



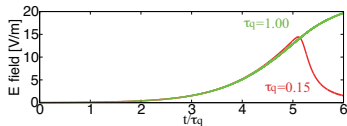
- Evolutions of current, loop voltage, and RE gen. rate on magnetic axis
- $\tau_q = 0.15$  ms ( $\tau_s^{ee}(3v_{th}) = 0.26$  ms)
- dotted: excluding the hot-tail
- solid: including the hot-tail

- The hot-tail effect enhances the RE current. (fig. a)
- the peak value of the primary RE gen. rate is an order magnitude greater. (fig. (c))
- The  $E$  field including the hot-tail sharply drops (fig. (b))
  - owing to the high primary RE gen. rate.
- Once the  $E$  field decreases, it re-increases gradually. (fig. (b))
  - The hollow  $E$  profile is filled by  $E$  diffusion.

# Hot-tail affects the RE distribution (JT-60U case)

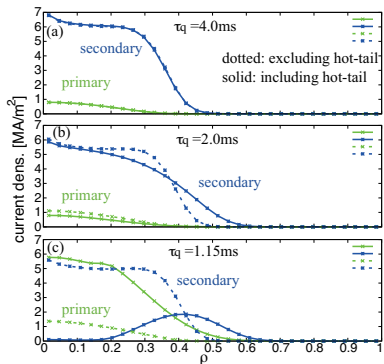


- left:  $\tau_q = 0.15$  ms (hot-tail effective)
- right:  $\tau_q = 1.0$  ms (ineffective)
- dashed curve: RE-nonRE boundary derived by Fussmann[NF 19, 327, (1979)]



- $E$  field is similar until  $t = 5\tau_q$  in both cases.
- The hot-tail effect makes  $f_e$  broad to the  $p_{\perp}$  direction.
- There are a lot of hot-tail electrons, which have finite perpendicular momentum (See  $t = 2 - 3\tau_q$ )
- Collisional pitch angle scattering also affects the broadening.

# RE current density profile (ITER like)

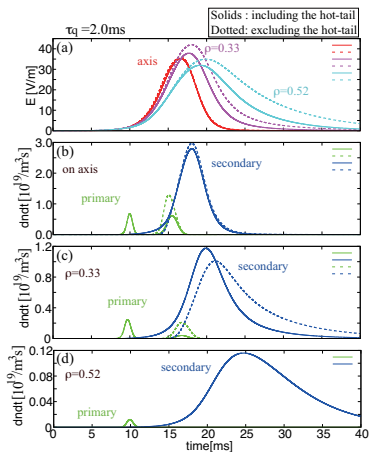


■  $\tau_s^{ee}(3v_{th}) = 3.88\text{ms}$

- RE  $j$  profiles are in good agreement for  $\tau_q < \tau_s(3v_{th})$ .
- For  $\tau_q = 2.0$  ms, the secondary RE  $j$  profile including the hot-tail effect becomes broader.
  - The hot-tail effect reduces the primary RE  $j$  around  $\rho = 0$  and enhances it in outer region ( $\rho > 0.4$ ) in invisible magnitude.
- For  $\tau_q = 1.15$  ms, the hot-tail effect dominates the primary RE  $j$  on-axis.
- The secondary RE  $j$  has a hollowed profile.



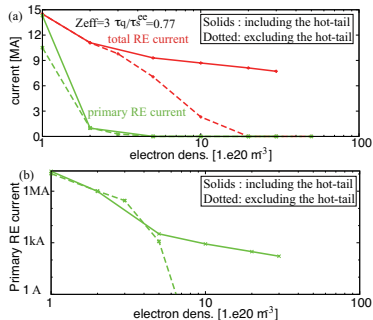
# Evolutions of $E$ and $dn/dt$ for $\tau_q = 2.0\text{ms}$ (ITER)



- (a):  $E$  field at several radial point.
- (b)-(d): RE gen. rate.

- The hot-tail effect divides the peak of the primary RE gen. into two. ( $t \sim 10$  and  $t \sim 15\text{ms}$ )
  - Former is the hot-tail generation.
- The hot-tail electrons can be REs even with the weak  $E$  ( $E/E_C \sim 20$  at  $t = 10\text{ms}$ ).
- REs generated at the earlier time trigger the secondary gen.
- Earlier onset on the secondary gen. maintains  $E$  weaker.
- Subsequent primary gen. decreases
  - The primary gen. is sensitive to  $E$ .

# Density dependence



- (a) Primary and secondary RE current against electron density.
  - (b) Focused on the primary RE
  - $\tau_q/\tau_s = 0.77$  (fixed)
  - Evolution of  $n_e$  is omitted.
- The high density  $n_e \sim 10^{21} \text{ m}^{-3}$  may suppress RE current (excluding hot-tail).
  - If  $\tau_q/\tau_s < 1$  is kept
    - Hot-tail maintains 0.1 – 1kA even if  $n_e > 10^{21} \text{ m}^{-3}$
    - Primary REs are multiplied to  $\sim \text{MA}$ .
  - Reliable  $\tau_q$  in mitigated plasma is required.

# Summary

- RE generation including hot-tail effect has been investigated.
- The hot-tail effect affects to the RE current, the evolution of  $E$ ,  $j$  profile, RE velocity distribution, and density dependence even in secondary dominant case.
- There are potentials that
  - RE velocity distribution affects to synchrotron rad.
  - If the mitigation shortens the thermal quench sufficiently, hot-tail effect makes innegligible RE current.