

Simulation study of electrostatic potentials produced by fast-ion population in toroidal plasmas

H. Yamaguchi* and S. Murakami

Department of Nuclear Engineering, Kyoto University

*JSPS fellow,

This work was supported by JSPS KAKENHI Grant Number 26420851 and 15J08296.



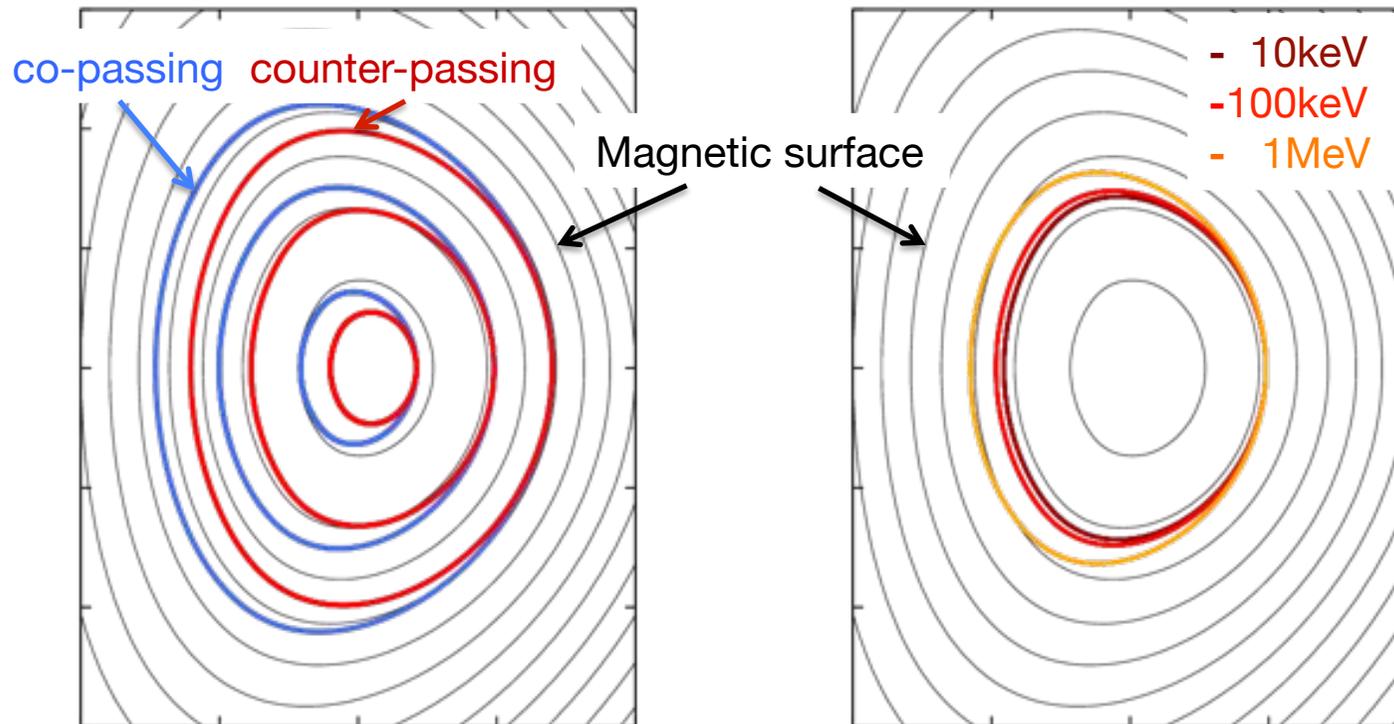
京都大学
KYOTO UNIVERSITY



Introduction

Fast ions in fusion plasmas

- Alpha particles produced by D-T fusion reaction ($E=3.5\text{MeV}$)
- Fast ions produced by NBI and ICRF heating ($E\sim 100\text{keV}-1\text{MeV}$)
- Primary heat source of plasma
- Potent heat load to the divertor => should be well-confined



Typical guiding-center orbit of co-passing and counter-passing fast ions in a tokamak plasma.

Energy dependence of guiding-center orbit of co-passing ions

Orbit of fast ions in toroidal plasmas

- Deviation from the magnetic surface is relatively large.
- Hardly trapped in the potential wells of the level of thermal energy
 - fast ions tend to be **non-uniformly** distributed over a magnetic surface.
 - may produce electrostatic (ES) potential varying on magnetic surface
 - Such fast-ion-induced ES potentials and their effect on plasma performance has not been investigated.

Objectives of this study

- We evaluate ES potentials produced by fast-ion non-uniformity in toroidal plasmas, on the basis of numerical simulations.
- We also consider **the presence of magnetic islands**, which may lead to **further localization of passing fast ions**.
- We investigate the effect of the ES potentials on fast ion confinement.

Fast ion model

GNET (Global NEoclassical Transport) code

$$\frac{\partial f_f}{\partial t} + (\mathbf{v}_{\parallel} + \mathbf{v}_D) \cdot \nabla f_f + \dot{\mathbf{v}} \cdot \nabla_{\mathbf{v}} f_f = C(f_f) + L_{\text{particle}}(f_f) + S_{\text{beam}}$$

f_f : fast-ion distribution function, \mathbf{v}_{\parallel} : parallel velocity, \mathbf{v}_D : drift velocity,
 C : Linear Coulomb collision operator, L_{particle} : particle loss term,
 S_{beam} : fast-ion source term (by HFREYA)

- We solve above equation for f_f in 5D phase space, using the GNET code [Murakami 2006, Nucl. Fusion] based on Monte Carlo technique.
- Guiding-center orbit is followed in Boozer coordinates with 6th-order Runge-Kutta-Hutta method.
- Pitch-angle and energy scatterings during energy slowing down
- Magnetic field and plasma geometry from VMEC
- **Extended to treat ES potentials with arbitrary Fourier modes.**

Electrostatic potential model

Adiabatic response of electrons

- ES force is assumed to be balanced by pressure gradient force in parallel direction (Boltzmann relation in parallel direction).

$$n_e = n_{e0} \exp\left(\frac{e\delta\Phi}{\kappa T_e}\right) \approx n_{e0} + \delta n_e \implies \delta\Phi = \frac{\kappa T_e}{e} \frac{\delta n_e}{n_{e0}}$$

- Equilibrium- and perturbed parts of electron density:

$$n_{e0} = n_{i0} + \langle n_f \rangle$$

$$\delta n_e = n_f - \langle n_f \rangle,$$

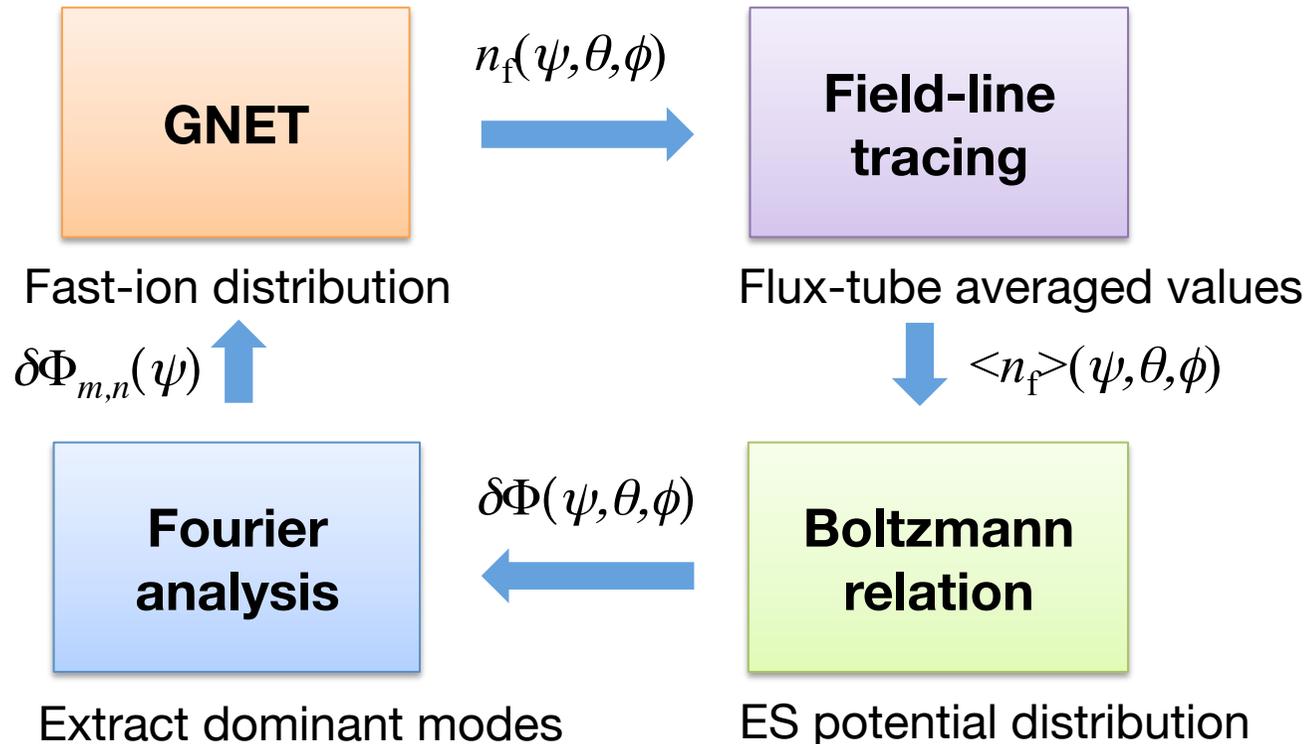
- T_e and n_{i0} are assumed to be constant along each field lines.
- $\langle n_f \rangle$: the flux-tube averaged fast-ion density (by field-line tracing)

$$\langle n_f \rangle \equiv \int n_f \frac{dl}{B} / \int \frac{dl}{B}$$

Iterative simulation

Data flow in the iteration

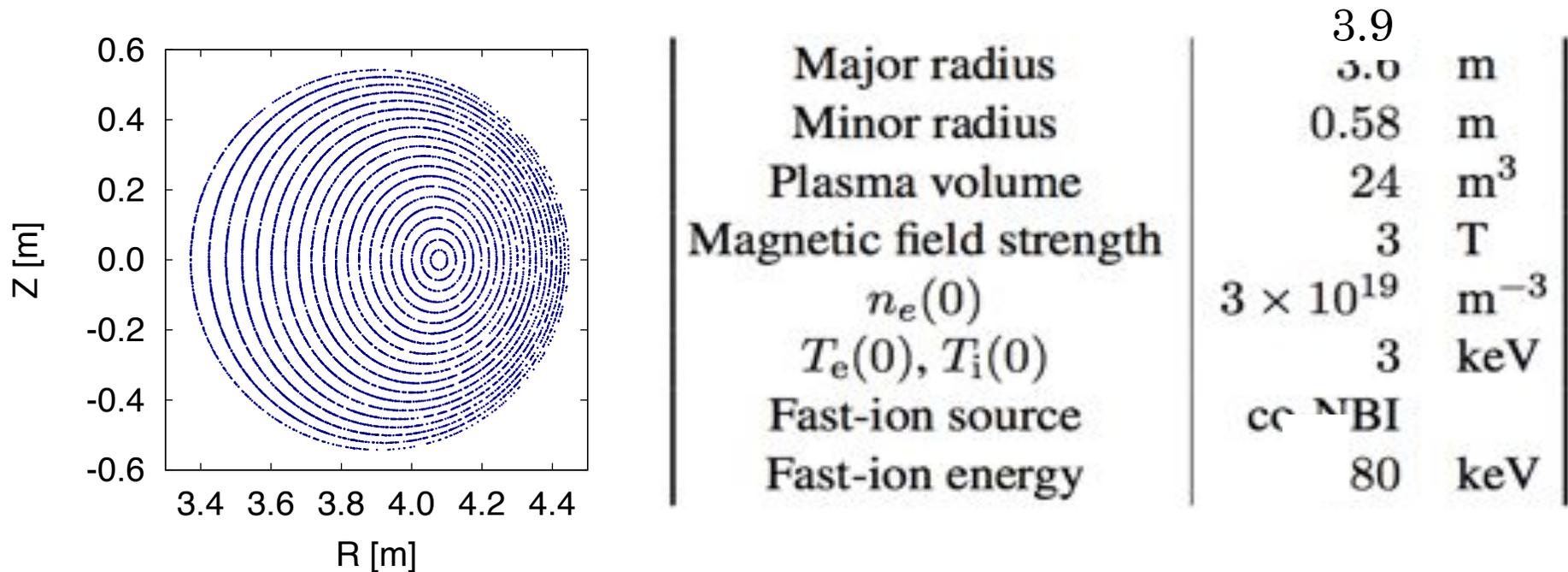
- We have extended GNET to **ES potentials with arbitrary Fourier mode**.
- Self-consistent fast-ion distributions and ES potentials can be obtained by the following iteration (background ion is assumed rest).



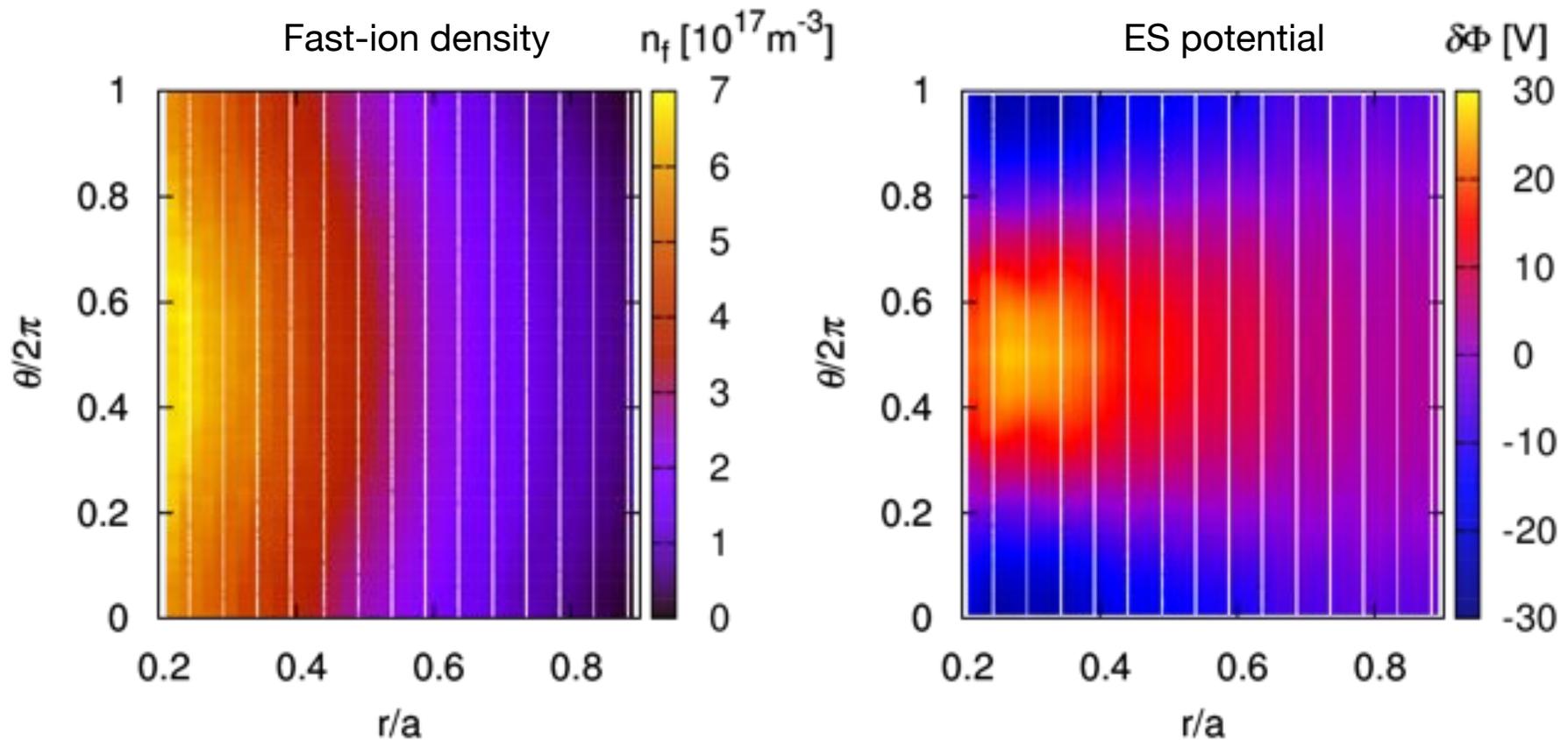
Simulation conditions

Assumed equilibrium, plasma profiles, and fast-ion source

- We consider a tokamak plasma with a circular cross-section.
- Electron density and electron- and ion temperatures at the center are $n_e=3 \times 10^{19} \text{m}^{-3}$ and $T_e=T_i=3 \text{ keV}$.
- A co-current NBI injected with $E=80 \text{ keV}$ is assumed. (no sub-components)
- NB absorption power is set to 5 MW.



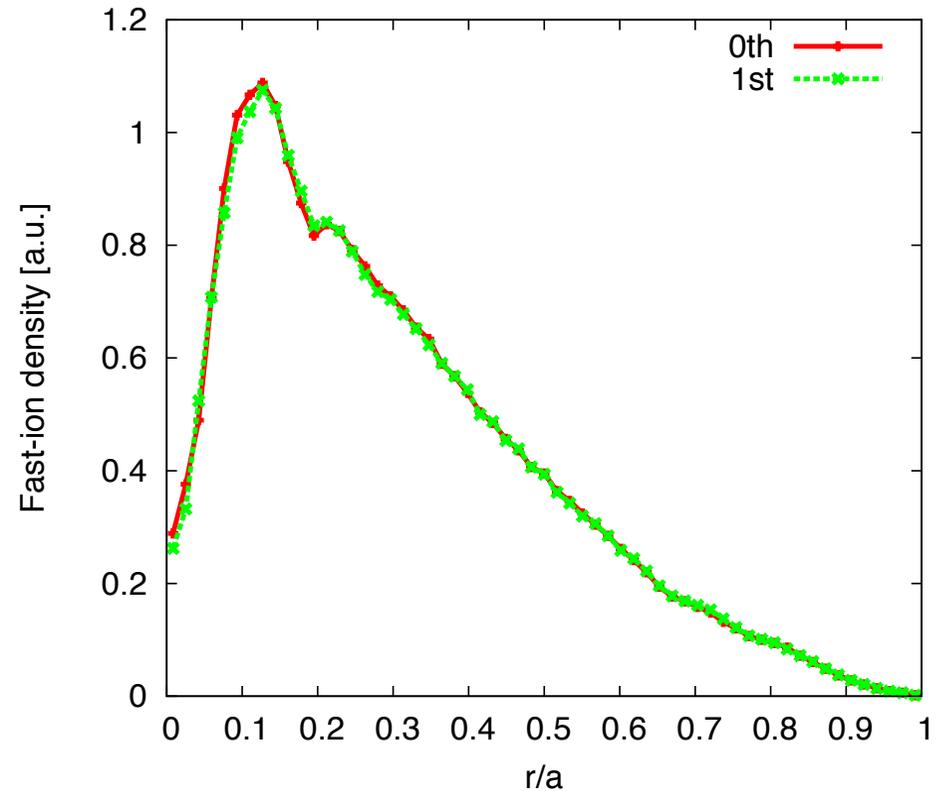
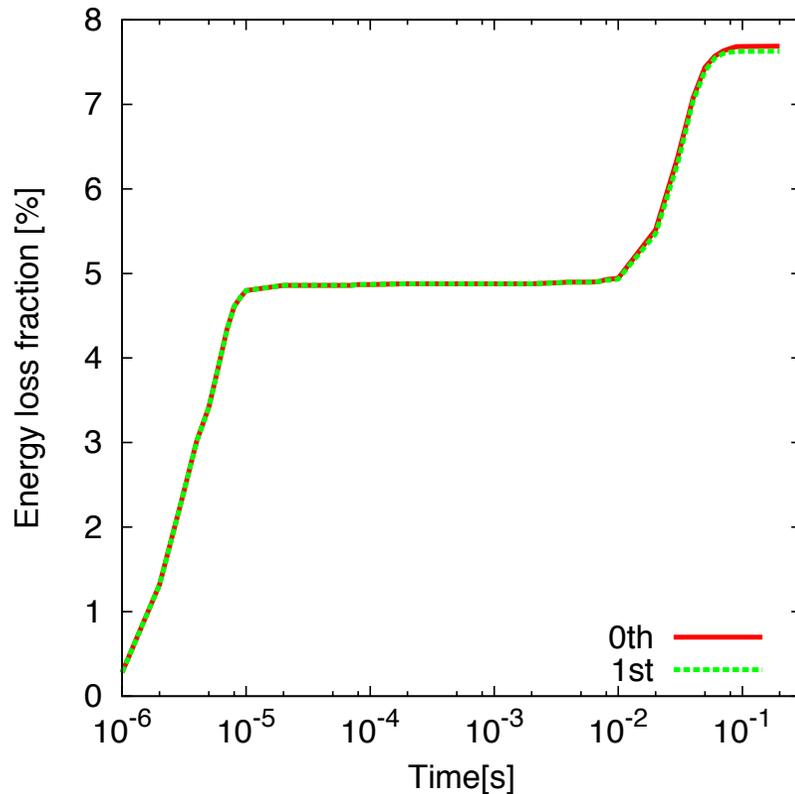
Simulation results I



- Fast-ion density varies over magnetic surfaces (white lines) with the dominant poloidal mode number of 1 due to toroidicity.
- ES potential $\sim 20\text{V}$ is produced by the non-uniformity of fast ions.

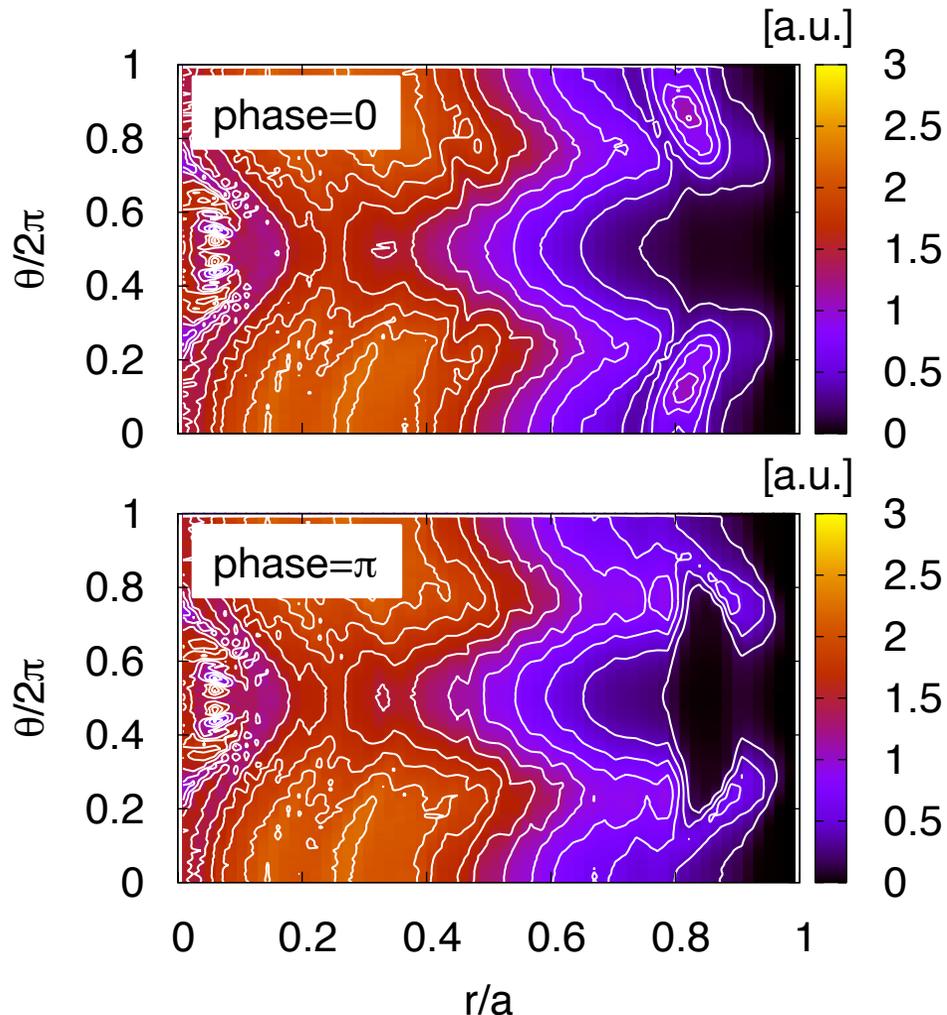
Result of iterative simulation

- In the absence of magnetic islands, no clear change was found before and after a single iteration.



Fast ion localization by magnetic islands

NBI beam pressure with magnetic island in the LHD



- Previously, we performed NBI heating simulation of LHD (Large Helical Device) plasma with magnetic islands by RMP (Resonant Magnetic Perturbation). (ITC25, 2015)
- We have found that the tangentially-injected fast beam ions form highly-localized beam ion pressure profiles near the resonant magnetic surface.
- Next, we investigate ES potentials produced by fast ions in presence of magnetic islands.

Magnetic island model

Perturbation model for magnetic island

- We use a well-known analytic form of magnetic perturbation, δB , producing magnetic island at resonant rational surface.

$$\mathbf{B}_{\text{total}} = \mathbf{B} + \delta \mathbf{B} \quad , \quad \delta \mathbf{B} = \nabla \times (\alpha \mathbf{B}) \quad , \quad \alpha = \alpha_{m,n}(r) \cos(m\theta - n\zeta + \delta)$$

$$\dot{\psi} = -\frac{1}{\Gamma} \left(\frac{e^2 B}{m} \rho_{\parallel}^2 + \mu \right) (g \partial_{\theta} B - I \partial_{\zeta} B) - \frac{1}{\Gamma} e (g \partial_{\theta} \Phi - I \partial_{\zeta} \Phi) + \frac{g_3}{\Gamma} \frac{e^2 B^2}{m} \rho_{\parallel}$$

$$\dot{\rho}_{\parallel} = -\frac{1}{\Gamma} \left(\frac{e^2 B}{m} \rho_{\parallel}^2 + \mu \right) (-g_2 \partial_{\theta} B + g_1 \partial_{\zeta} B) - \frac{1}{\Gamma} e (-g_2 \partial_{\theta} \Phi + g_1 \partial_{\zeta} \Phi) \\ - \frac{g_3}{\Gamma} \left(\frac{e^2 B}{m} \rho_{\parallel}^2 + \mu \right) B'$$

$$\dot{\theta} = \frac{g}{\Gamma} \left\{ \left(\frac{e^2 B}{m} \rho_{\parallel}^2 + \mu \right) B' + e \Phi' \right\} - \frac{g_2}{\Gamma} \frac{e^2 B^2}{m} \rho_{\parallel}$$

$$\dot{\zeta} = \frac{-I}{\Gamma} \left\{ \left(\frac{e^2 B}{m} \rho_{\parallel}^2 + \mu \right) B' + e \Phi' \right\} + \frac{g_1}{\Gamma} \frac{e^2 B^2}{m} \rho_{\parallel}$$

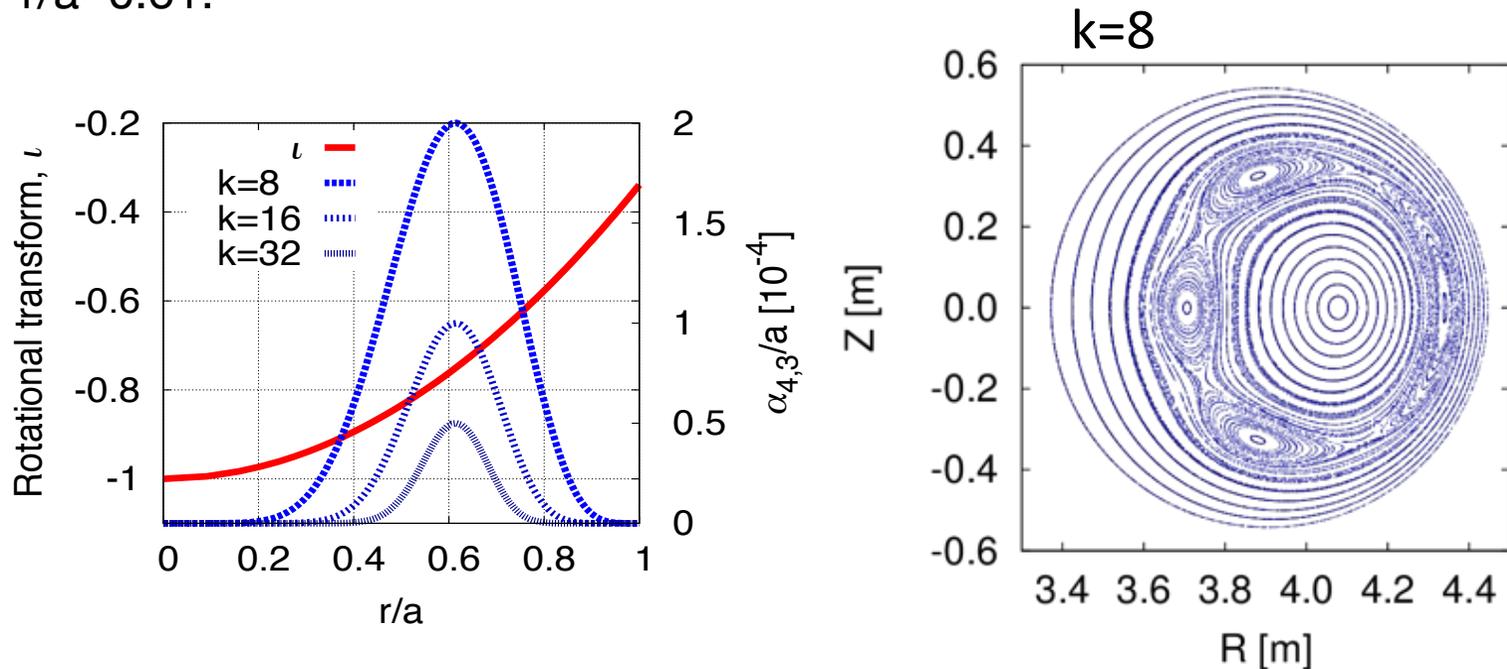
$$\Gamma = e \left[g \{ (\rho_{\parallel} + \alpha) I' + \alpha' I + 1 \} - I \{ (\rho_{\parallel} + \alpha) g' + \alpha' g - \iota \} \right]$$

$$g_1 = (\rho_{\parallel} + \alpha) I' + \alpha' I + 1$$

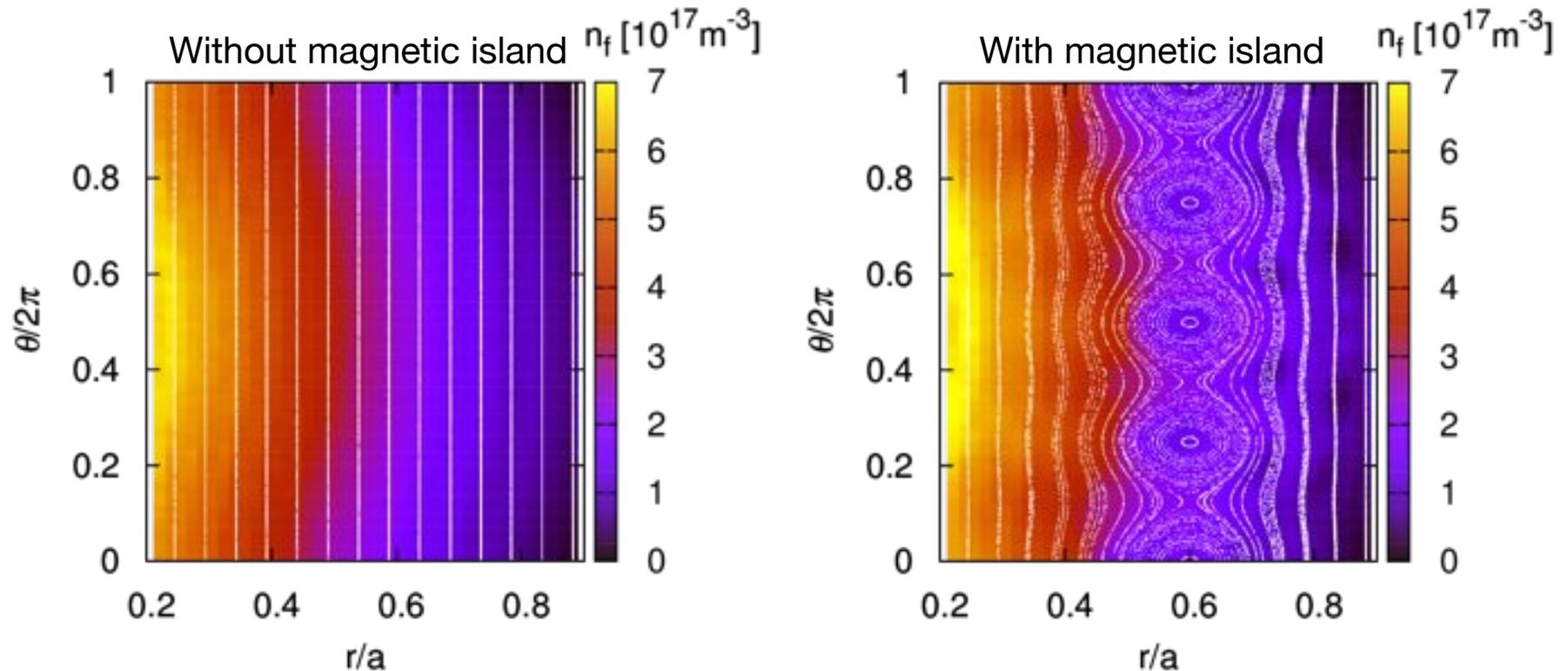
$$g_2 = (\rho_{\parallel} + \alpha) g' + \alpha' g - \iota$$

$$g_3 \equiv g \frac{\partial \alpha}{\partial \theta} - I \frac{\partial \alpha}{\partial \zeta},$$

- We consider a static (m,n)=(4,3) mode whose resonant surface locates at $r/a \sim 0.61$.

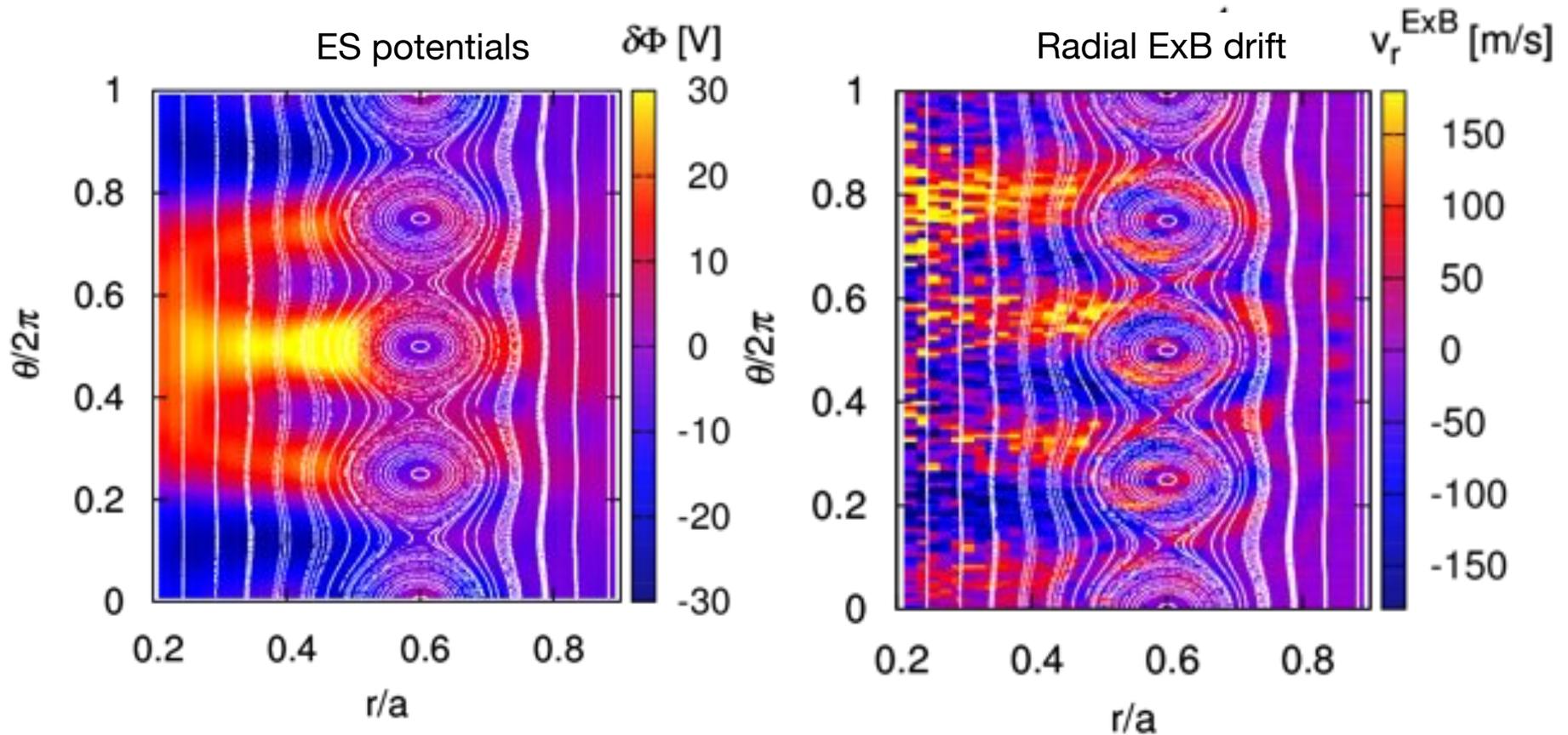


Fast ion density with magnetic island



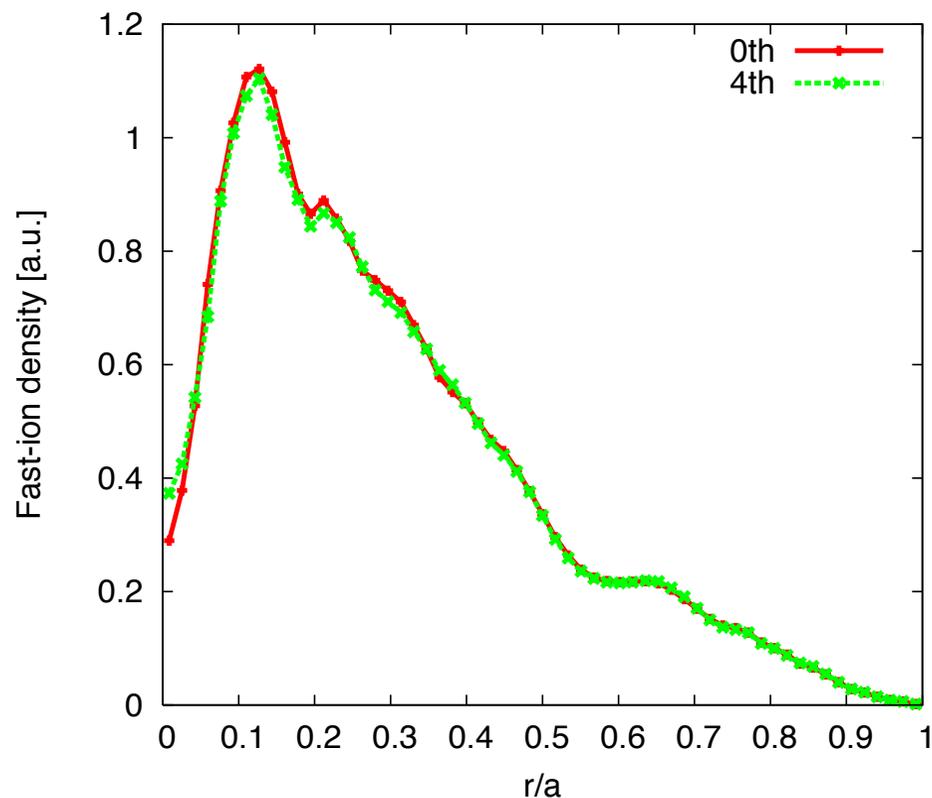
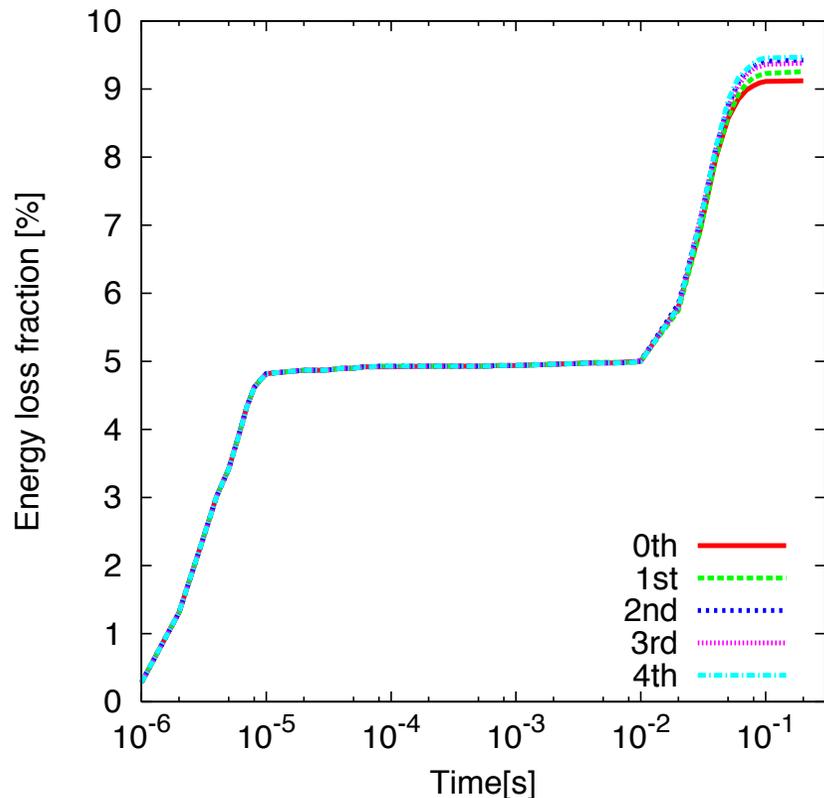
- Equilibrium magnetic surfaces are distorted and torn into magnetic island at the resonant rational surface by the superimposed magnetic perturbation.
- Drift island structure appears in the fast-ion density profile in the presence of magnetic island.

ES potentials with magnetic island



- Potential wells are formed along the axis of drift islands.
- ExB convection cells across magnetic surfaces are formed.

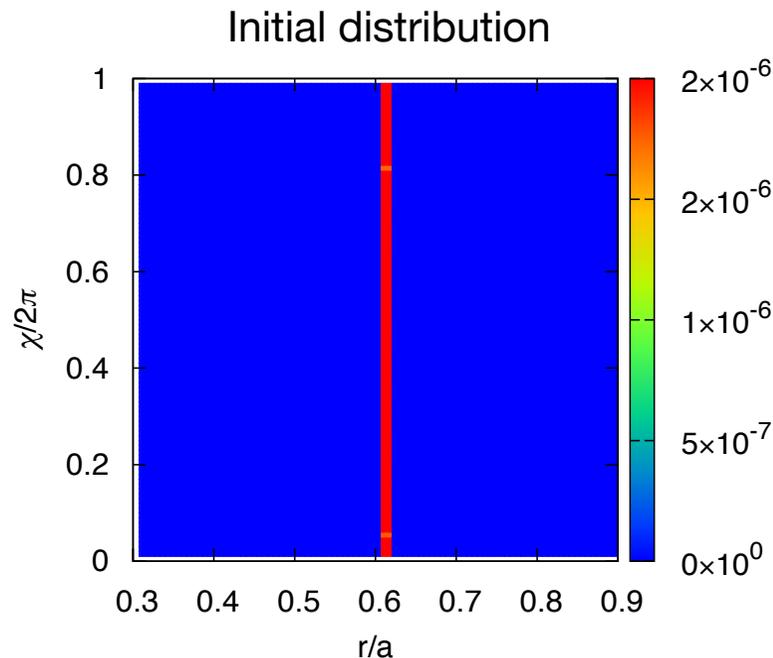
Effect on fast ion transport



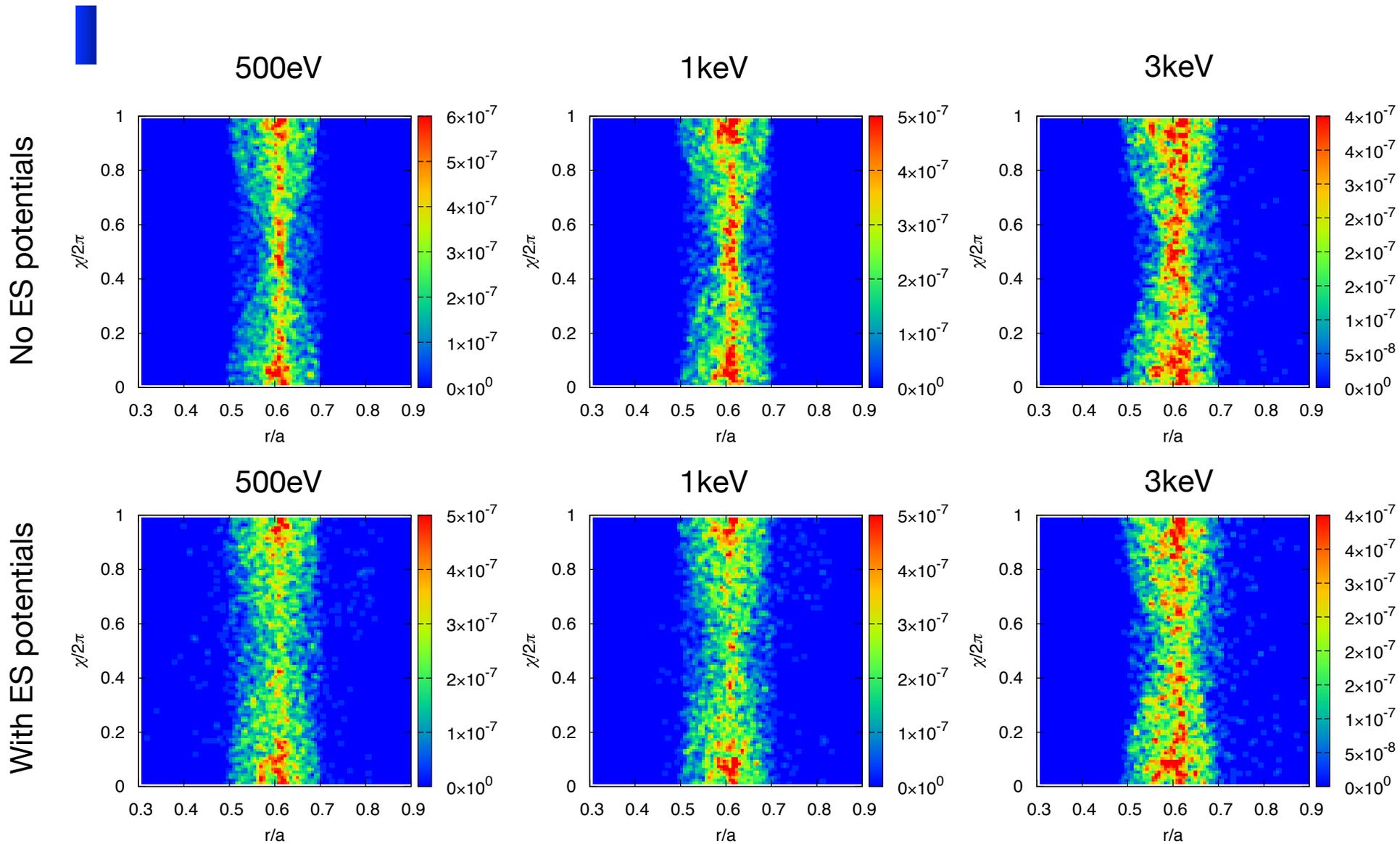
- We performed several iterative simulation of NBI fast ions and ES potentials.
- The energy loss fraction of fast ions slightly ($<1\%$) increased and the fast ion density decreased in the central region after 4 iterations.
- The effect on fast ion transport is very small.

Effect on thermal ion transport

- We distributed 5000 isotropic and mono-energetic test particles initially at the resonant surface.
- We followed the test particles with three different energy (500eV, 1keV, and 3keV) for 10 ms and investigated the effect of fast-ion-induced ES potentials on spatial diffusion in presence of magnetic islands.



- $\chi = 4\theta - 3\phi$
- In the absence of ES potentials, ions tend to remain on the resonant magnetic surface.
- We can see clear enhancement of diffusion due to ES potentials especially near the X point for each energy.



Summary

- We have evaluated electrostatic (ES) potentials produced by non-uniformity of fast-ion density along field lines in toroidal plasmas, using GNET code.
- In a tokamak plasma, an ES potential ~ 30 V with the dominant toroidal mode number of 1 is formed due to toroidicity.
- Magnetic islands have strong localization effect on fast ions
- The effect of ES potential in the presence of magnetic island
 - On fast ions: very small (slight decrease in the central density)
 - On thermal ions: clear change in the spatial diffusion was found.

Future task

- Detailed study on the particle diffusion (evaluation of diffusion coefficient)
- Dependency of fast-ion-induced ES potentials on the island phase
- ES potentials formation by NBI fast ions in the LHD and its effect on particle transport
- Diffusion/confinement of fast ions in low-frequency, rotating magnetic island of NTM