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原子力機構におけるトカマク統合 コードの開発と最近の研究成果

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Introduction

Issues of fusion burning plasmas: Control of plasmas with strong complexity and autonomy

- Multi-scale Physics
 - Consit of the wide spacial scale and the time scale
 - Consist of complexibility and self-organization
- Recurring Process
 - Bulk plasma determines a-heating profile, α -heating profile determines the bulk plasma profile
 - Nonlinearity: multi equilibrium point, bifurcation, discontinuity
- Steady-state control
 - α -heating>external heating: Reduction of heating power for the control
 - Bootstrap current>externally driven current: Reduction of CD and momentum input for the control
 - Existence of the steady state solution, controllability, stiffness

Background

- To progress the research of burning plasma, it is necessary to understand complex features of the advanced tokamak plasmas.
 - Because, in the steady state, plasma has very wide time scale and spatial scale, and it has complex physics such as turbulence, transport, MHD, wave-particle interaction, plasma-wall interaction, atomic and molecular physics, and so on



- Exploration and understand of the complex plasma are big issues.
 - High confinement, high beta, high bootstrap fraction, high radiation fraction, control of ash and impurity, α -heating, etc.
 - Control of the complex plasma are also serious issues.
 - Strong coupling of each physics mechanism, i.e. autonomous plasmas

Modeling and integration of models are useful means for the understand and the prediction of the burning plasma.

Strategy of Modeling / Integration of model

To make the integrated model, validation is necessary based on the fundamental researches of experiment and simulation.



Core plasma	Core MHD stability: Sawtooth, NT Heat/particle transport: ITB Current drive Wave-particle interaction α-Heating, Energetic particles	M TOPICS-IB: TOPICS extended to the Integrated
Edge/pedestal	Pedestal transport: ETB Edge MHD stability: ELM SOL transport, recycling	modeling for Burning plasma
SOL/Divertor	SOL/Divertor plasma transport Neutral particles Impurity particles Interaction with wall, A & M	SONIC: SOLDOR/ NEUT2D/ IMPMC
Momentum	Plasma rotations Radial electric field J×B torque	TASK:
Wave	RF Heating: IC Alfvén Eigenmode: TAE, RSAE Fokker-Planck	WM/TX

核燃焼プラズマ統合コード開発構想

BPSI: Burning Plasma Simulation Initiative

大学, 核融合研, 原子力機構等の研究協力





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TOPICS-IB (TOPICS extended to Integrated simulation for Burning plasmas)

TOkamak Prediction and Interpretation Code System

- 1D transport & 2D MHD equilibrium
- Time-dependent / Steady-state analysis of JT-60U experiment
- Simulation with transport model
 - Neoclassical : MI method or NCLASS, Anomalous : CDBM, GLF23, MMM95
- **Components and Related codes**
- NB & High energy particle : 1D or 2D FP, 3D Monte-Carlo (OFMC)
- EC : Ray tracing & Relativistic FP (EC-Hamamatsu)
- MHD : Kink / Ballooning / Peeling (MARG2D)
- Impurity : 1D transport (IMPACT), 2D Monte-Carlo (IMPMC)
- Neutral : 2D Monte-Carlo
- Radiation : Synchrotron (CYTRAN)
- SOL / Div. : Five-point model (D5PM), 2D Fluid & Monte-Carlo (SONIC)

Recent works

- Core-SOL-div. integration for ELM study (loss & cycle) (Hayashi, IAEA08)
- Integration with 1d1v FP code of α particles (Ozeki, Hamamatsu, JSPF08)

周辺部局在モードによるエネルギー損失

背景 核融合炉への課題:プラズマ閉込め容器への熱負荷の低減





ELMによるエネルギー損失: 多装置の実験データは、プラズマの衝突周波数に依存 ITERの高温(低衝突)プラズマで過大な損失?

衝突周波数依存性の発生機構が未解明の課題

Simulation of ELM energy loss and cycle



Integrated modeling reproduces the collisionality dependence of ELM energy loss.



ELM energy loss in the follwoing ELMs deviates a little from that at first ELM, but the collisionality dependence is almost the same.

As found in the first ELM (H-modeWS07), the following physics cause the dependence.

- Electron : Bootstrap current broadens the region of ELM enhanced transport and SOL parallel conductive transport decreases the SOL temperature in the low collisionality. (NF07)

- Ion : $T_i > T_e$ in the low collisionality due to the ineffectiveness of equipartition, which enhances ion convective and CX losses.

Reduction of total energy loss is comparable with that in experiments.

Integrated Simulation of ELM Energy Loss and Cycle in Improved H-mode Plasmas (Hayashi, IAEA08)

- Integrated code TOPICS-IB clarified that steep pressure gradient inside the pedestal top broadens the region of ELM enhanced transport and enhances the ELM energy loss.

- Transport model of **pedestal neoclassical transport connected to SOL parallel transport** reproduces the **experimentally observed collisionality dependence of inter-ELM transport**. Inter-ELM energy confinement time agrees with JT-60U scaling.



Energetic particle loss by AE modes



[[]M.Ishikawa, IAEA, FEC, 2006]

AE modes in a weak shear plasma





 (A) Sn reduced in 20~30% when AE modes existed.
 Degradation of high energy particles

(B) Measured Sn was close to the OFMC calculation: ~ classical confinement

Reduction rate of Sn increased in AE mode flucuation

Plausible candidate is MHD resonance particle loss.

Simulation of Autonomous and Recurrence Process



 $f(v,\rho,t)$: Velocity distribution function of alphas v: velocity, ρ : minor radius $\frac{\partial}{\partial t}f(v,\rho,t) + \frac{1}{\rho}\frac{\partial}{\partial \rho}\left|\rho\left(-D(v,\rho)\frac{\partial f}{\partial \rho} + V_{AN}(v,\rho)f\right)\right| = \sum_{j}C_{j}(f) + S(v,\rho) - L(f)$ $\underline{D(v,\rho)} = D_{NC}(v,\rho) + D_{AN}(v,\rho)$ **2D Fokker-Planck Equation** $\bigcirc D_{_{NC}}(v,
ho)$: Neo-classical like diffusion $D_{NC}(v,\rho) = \sqrt{\frac{2\varepsilon}{1+\varepsilon}} \frac{\Delta_b^2}{\tau}, \Delta_b = \frac{2v\cos\alpha}{\Omega}, \tau_{eff} = \frac{4\varepsilon}{1+\varepsilon} \tau_{\perp}(v), \varepsilon = \frac{\rho}{R_{\bullet}}, \tau_{\perp}(v) \quad : \text{Pitch angle scat. time}$ $\bigcirc \ D_{_{\!\!A\!N}}(v,
ho)\,\&\,V_{_{\!\!A\!N}}(v,
ho)$: Anomalous diffusion & convection model of the resonance with MHD instability $\bigcirc C_i(f)$: Collision term colliding with *j*-th bulk plasma $n_s(\rho)$ and $T_s(\rho)$ are calculated by TOPICS-IB \bigcirc $S(v, \rho)$: Particle source by fusion reaction and NBI L(f) : Loss term removes slowed down particles $\bigcirc L(f) = \frac{f(v,\rho,t)}{\tau_{at}(\rho)} \exp\left(-\frac{m_{\alpha}v^{2}}{\lambda T_{ach}(\rho)}\right) \qquad \lambda \text{ is constant.} \Rightarrow \lambda = 3 \text{ in this case} \\ \tau_{th}(\rho) \text{ : thermal collison time}$

Simulation of F-P equation for fixed background

- Effects of the anomalous transport on the α -particle pressure/heating
- Assumption: : neo-classical like α -diffusion, the anomalous by the MHD fluctuation and zero flow velocity



•Time trace and profiles for the simulation with (red) and without(green) anomalous

On of anomalous: $\frac{\partial p_{\alpha}}{\partial \rho}$

 $\left\| \sum_{max} \geq \frac{\partial p_{\alpha}}{\partial \rho} \right\|_{on}$

Off of anomalous:

 ∂p_{α}

Integrated simulation results

Critical gradient of α pressure

HFACTX

HFACTX



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Integrated Divertor Code for Fusion Reactor

To investigate the power and particle control in tokamak reactor, integrated divertor codes have been developed.



Integrated Divertor Code in JAEA



IMPMC uses (1) Diffusion model using Langevin analytical solution to extend time step, (2) MPI with many particles to reduce MC noise, (3) Particle reduction scheme to obtain steady state.

Simulation of X-point MARFE with SONIC

To reduce the high heat load onto the divertor plate, the control method for impurity retention in the divertor region should be established.



Hydrocarbons sputtered from the dome contribute to the enhanced radiation near the X-point.

Plan of integration



EDDY/IMPMC Simulation

for Methane Breakup Model



Comparison with Simple Methane Breakup





In detached plasma, simple model is relatively good approximation

In attached plasma, the methane breakup can not be simplified

The dome with a small sticking coefficient enhances the contam- ination of carbon into the main plasma.

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M. Honda and A. Fukuyama, J. Comp. Phys. 227 (2008) 2808



Toroidal rotation induced by charge separation (

- Ionization of fast neutrals from NB produces electrons and fast ions. 1)
- Owing to the toroidal drift, the trapped ion especially deviates from the birth flux surface. 2)
- The charge separation continues as long as NB is injected, causing local charge imbalance. 3)

30°

(b) 30°

(a) 0°

3.8

R (m)

(b)

- A radial current flows in the bulk plasma to maintain quasi-neutrality.
- A resultant $j \times B$ torque drives the toroidal rotation. 5)

TASK/TX simulations with OFMC

important

- have reproduced the toroidal rotation by charge separation due to near perpendicular NBI.
- have clarified that horizontally-injected NB (case (a)) drives the \geq rotation most efficiently.

Banana width and magnetic field at birth position



Effect of ripple loss of fast ions on toroidal rotation in JT-60U



- In JT-60U experiments, effects of ripple loss of fast ions on toroidal rotation have been extensively studied [5].
 - Installation of ferritic steel tiles (FSTs) reduces the toroidal field ripple amplitude roughly in half.
 - The discharge after installing FSTs (smaller ripple loss) was compared with the otherwise equivalent case with the ripple loss amplitude before installing FSTs.

Loss of fast ions due to the ripples drives the plasma rotation in the counter direction.

From the experimental analysis, the counter rotation rather than the pressure gradient has strong ties with the formation of the inward E_r .

[5] M. Yoshida, et al., PPCF 48 (2006) 1673



Ripple loss channels



- In our model ripple transport is assumed to be governed by three different types of transport mechanisms.
 - Vertical drift (collisionless ripple-trapping) loss

Beam ions trapped in ripple well drift vertically with grad B drift velocity v_{d} .

$$u_b^{rp} \approx \frac{\Delta r}{z/v_d} = \frac{z}{r}v_d = \frac{m_b v_b^2}{2Z_b eBR} \frac{NqR\delta}{r}, \text{ where } \Delta r = \frac{z^2}{r}, z = NqR\delta$$

- Collisional diffusive loss

High collisionality regime (ripple-plateau diffusion) [7]

$$D_{RP} \cong v_b \Delta r^2$$
, where $v_b \cong \frac{\varepsilon^{1/2} v_b}{4\sqrt{2}K(\xi)qR} \cong \frac{\varepsilon^{1/2} v_b}{10.5qR}$, $\Delta r^2 = \pi N \left(\frac{\varepsilon}{q}\right)^3 \delta^2 \rho^2$

Low collisionality regime (ripple-resonance diffusion) [8]

$$D_{CB} \cong \frac{N^{9/4} q^{13/4} R \rho_b \delta^{3/2}}{\varepsilon^{5/2}} v_D$$

- Collisionless stochastic diffusion [9]

Diffusion coefficient is the same as ripple-plateau diffusion coefficient, D_{RP} . Stochastic diffusion occurs where the following GWB criterion is satisfied.

 $\delta_{c} \leq \left[\left(\frac{\pi N q}{\epsilon} \right)^{3/2} \rho \frac{dq}{dr} \right]^{-1} \begin{bmatrix} 7 \end{bmatrix} \text{A. H. Boozer, PoF 23 (1980) 2283} \\ \begin{bmatrix} 8 \end{bmatrix} \text{V. Ya. Goloborod'ko et al., Physica Scripta T16 (1987) 46} \\ \begin{bmatrix} 9 \end{bmatrix} \text{R. J. Goldston et al., PRL 47 (1981) 647} \end{bmatrix}$





Summary

Core integrated code TOPICS-IB

- Core-SOL-div. integration for study of ELM loss & cycle

- Integration with 1d1v FP code for study of AE modes effect on α particle confinement

Divertor integrated code SONIC

- IMPMC-EDDY integration for study of methane breakup effect on X-point MARFE

Two-fluid core code TASK/TX

- TASK/TX -OFMC integration for study of toroidal rotation induced by charge separation of fast ions (due to ripple & perpendicular NBI).